

---

# Linear Systems and Mixed Finite Elements in Hydrogeology

Hussein Hoteit<sup>†\*</sup>, B. Philippe<sup>†</sup>, J. Erhel<sup>†</sup>  
R. Mose<sup>\*</sup> and Ph. Ackerer<sup>\*</sup>

<sup>†</sup>IRISA-INRIA, Rennes

<sup>\*</sup>Institut de Mécanique des Fluides et des Solides,  
Strasbourg

Sparse Days at Cerfacs

June 24-25, 2002

---

## Outline

- ▷ **Mixed and Mixed-Hybrid finite element approximations.**
  - ⇒ model flow problem in porous media.
  - ⇒ variational formulations and the resulting linear systems.
- ▷ **Effect of the mesh geometry on both approximations.**
  - ⇒ flat elements and condition numbers.
  - ⇒ symbolic and numerical computing (stability analysis).
- ▷ **Behavior of both methods on heterogeneous media.**
  - ⇒ condition number estimates of the algebraic linear systems.
  - ⇒ numerical error estimates of the computed velocity.
  - ⇒ choice of linear system solvers and preconditioning.
- ▷ **Some numerical experiments to validate the theoretical results.**

The governing transient equations for the pressure head  $d$  and Darcy's velocity are given by:

$$\begin{aligned}
 & \frac{\partial d}{\partial t} + \Delta \cdot n = f && \text{in } \Omega \times (0, T), \\
 & -\mathcal{K} \Delta d = n && \text{in } \Omega \times (0, T), \\
 & d = d_0 && \text{in } \Omega, \\
 & d = d_D && \text{on } \Gamma_D \times (0, T), \\
 & n \cdot \nu = b_N && \text{on } \Gamma_N \times (0, T).
 \end{aligned}
 \tag{1}$$

where,

$f$  : sink/source term;

$\mathcal{K}$  : tensor of hydraulic conductivity;

$s$  : storage coefficient;

$\nu$  : outward unit normal vector.

## Mixed and Mixed-Hybrid finite element methods

- The mass is conserved locally (over each element  $K$ )

$$\int_{\partial K} \nu \cdot n = \int_K \Delta \cdot n = \sum_{\partial K} \frac{\partial}{\partial x} \cdot flux.$$

- The pressure head and its gradient are approximated simultaneously with the same order of convergence

- They can easily handle full tensors of permeability

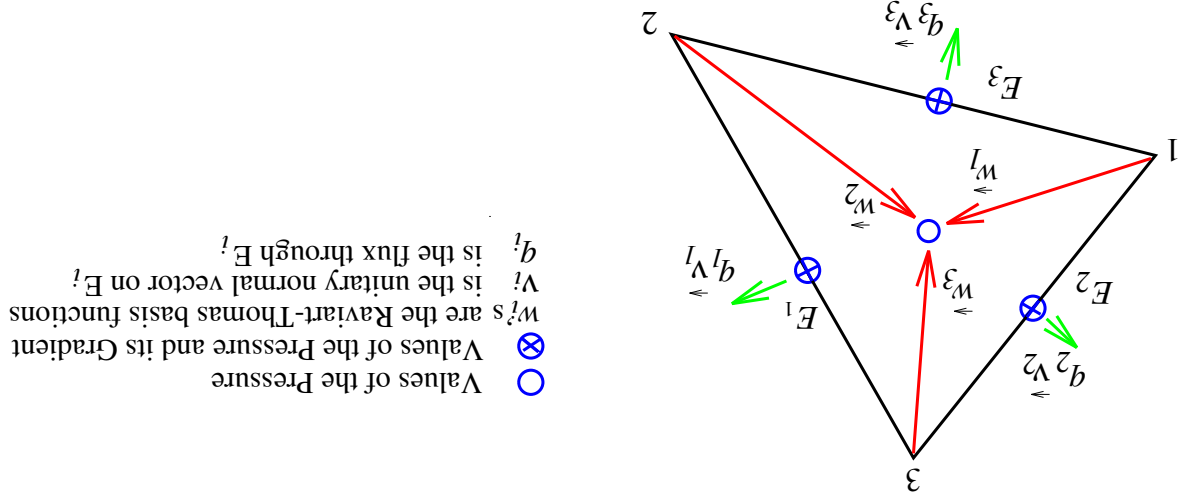
- The hybridization technique aims to enforce the continuity of the fluxes across the interelement boundaries

$$\int_E \nu \cdot n_{K_1} = \int_E \nu \cdot n_{K_2}, \quad E = \partial K_1 \cap \partial K_2.$$

The basis functions of the 3-dimensional space  $RT_0(K)$  are defined as

$$w_{K,E_i} = \frac{1}{2|K|} \begin{pmatrix} x - x_i \\ y - y_i \end{pmatrix} \quad i = 1, \dots, 3, \quad (2)$$

where  $|K|$  denotes the measure of the triangular element  $K$ .



⊙ Values of the Pressure and its Gradient  
 ⊗  $w_i$ 's are the Raviart-Thomas basis functions  
 $v_i$  is the unitary normal vector on  $E_i$   
 $q_i$  is the flux through  $E_i$

Nodal points and basis functions on triangular elements.

–  $p_K$ : the mean of the pressure in  $K$ .

–  $tp_{K,E}$ : the mean of the pressure on  $E$ .

The MFE formulation reads as :

Find  $(u_h, p_h) \in RT_0^{q,N}(T) \times M_0(T)$ , such that

$$\left\{ \begin{aligned} \int_{\Omega} (\mathcal{K}^{-1} u_h) \cdot \chi_h dx + \int_{\partial\Omega} p_D \nu \cdot \chi_h d\ell &= \int_{\Omega} p_h \Delta \cdot \chi_h dx \quad \forall \chi_h \in RT_0^{q,N}, \\ \int_{\Omega} s \frac{\partial p_h}{\partial t} \phi_h dx + \int_{\Omega} \Delta \cdot u_h \phi_h dx &= \int_{\Omega} f \phi_h dx \quad \forall \phi_h \in M_0. \end{aligned} \right.$$

(3)

The MHFE formulation reads as :

Find  $(u_h, p_h, tp_h) \in RT_0(T) \times M_0(T) \times \mathcal{N}_0^{p,D}(E)$  such that

$$\left\{ \begin{aligned} \int_{\Omega} (\mathcal{K}^{-1} u_h) \cdot \chi_h dx + \int_{\partial K} tp_h \nu_K \cdot \chi_h d\ell &= \int_{\Omega} p_h \Delta \cdot \chi_h dx \quad \forall \chi_h \in RT_0, \\ \int_{\Omega} s \frac{\partial p_h}{\partial t} \phi_h dx + \int_{\Omega} \Delta \cdot u_h \phi_h dx &= \int_{\Omega} f \phi_h dx \quad \forall \phi_h \in M_0, \\ \sum_{K \in \mathcal{T}_h} \int_{\partial K} u_h \nu_K d\ell &= \int_{\partial\Omega} u_h \nu_N d\ell \quad \forall \chi_h \in \mathcal{N}_0^{p,D}. \end{aligned} \right.$$

(4)

By using  $w_{K,E}$  as test functions in (3), we get

$$B^K Q_K = p^K e - T_P^K \quad K \in \mathcal{T}_h, \quad (5)$$

where,

$B^K$  : is a  $3 \times 3$  symmetric positive definite matrix (elementary matrix) (trix)

$$(B^K)_{E,E'} = \int_K w_{K,E}^T K_K^{-1} w_{K,E'} dx.$$

$p^K$  : the pressure average over  $K$ .

$e$  : vector of dimension 3 with unitary entries.

$T_P^K$  : vector of pressure averages on  $\partial K$ .

$Q_K$  : the fluxes across  $\partial K$ .

By inverting  $B_K$ , it is possible to eliminate the flux unknown. Thus, we obtain :

$$(9) \quad R^T P - M^T p + V = 0$$

These matrices are defined as follows :

$R^T$  is a sparse matrix of dimension  $N_{\mathcal{E}} \times N_T$

$$R = [R_{K,E}]_{N_T, N_{\mathcal{E}}}, R_{K,E} = \alpha_{K,E} = \sum_{E' \subset \partial K} (B_{-1}^k)_{E,E'}$$

$M$  : is a  $N_{\mathcal{E}} \times N_{\mathcal{E}}$  sparse matrix

$$M = [M_{E,E'}]_{N_{\mathcal{E}}, N_{\mathcal{E}}}, M_{E,E'} = \sum_{\partial K \supset E, E'} (B_{-1}^k)_{E,E'}$$

Integrating (4) leads to :

$$(7) \quad s_K |K| \frac{\partial p_K}{\partial t} + \sum_{E \subset K} q_{K,E} = f_K \quad K \in \mathcal{T}_h,$$

In matrix form, (7) is rewritten :

$$(8) \quad S \frac{dP}{dt} + DP - RTP = F,$$

where

$$S = [S_{K,K}]_{N_{\tilde{\sigma}}, N_{\tilde{\sigma}}} \text{ diagonal matrix:}$$

$$S_{K,K} = |K| s_K$$

$$D = [D_{K,K}]_{N_{\tilde{\sigma}}, N_{\tilde{\sigma}}} \text{ diagonal matrix:}$$

$$D_{K,K} = \alpha_K = \sum_{E \subset \partial K} \alpha_{K,E};$$

The linear systems derived from the MFE

The **MHFE** formulation gives rise a linear system of unknowns  $P$  and  $T_p$

$$(9) \quad \begin{pmatrix} \frac{dP}{dt} \\ \frac{dT_p}{dt} \end{pmatrix} + \overbrace{\begin{pmatrix} D & -R^T \\ -R & M \end{pmatrix}}^{\mathcal{J}} \begin{pmatrix} T_p \\ P \end{pmatrix} = \begin{pmatrix} S \\ 0 \end{pmatrix} + \begin{pmatrix} 0 \\ V \end{pmatrix}.$$

By discretizing (9) by Backward Euler scheme, the problem is reduced to solve

$$(M - \Delta t N) T_p^n = R^T G_{n-1}^T (S P_{n-1} + \Delta t F_n) + V_n.$$

where  $G = (S + \Delta t D)$ ,  $N = R^T G^{-1} R$ .

$A = (M - \Delta t N)$  can be seen as a Schur complement in S.P.D matrix.

The linear systems derived from the MFE

The MFE formulation gives rise a linear system of unknowns  $P$  and  $Q$

$$(10) \quad \begin{pmatrix} \frac{dP}{dt} \\ \frac{dQ}{dt} \end{pmatrix} = \overbrace{\begin{pmatrix} 0 & -\widetilde{R}^T \\ \widetilde{R} & 0 \end{pmatrix}}^{\mathcal{L}} \begin{pmatrix} Q \\ P \end{pmatrix} = \begin{pmatrix} \widetilde{V} \\ \widetilde{F} \end{pmatrix}.$$

By discretizing (10) by Backward Euler scheme, the problem is reduced to solve

$$(\widetilde{M} + \Delta t \widetilde{N}) Q_n = \widetilde{R}^T (P_{n-1} + \Delta t S^{-1} \widetilde{F}_n) + \widetilde{V}_n.$$

where  $\widetilde{N} = \widetilde{R}^T S^{-1} \widetilde{R}$ .

$\widetilde{A} = (\widetilde{M} + \Delta t \widetilde{N})$  can be seen as a Schur complement in S.P.D matrix.

## Effect of the spatial discretization

The linear systems derived from the MFE

Inverting method of the elementary matrix  $B_K$

This matrix is defined by

$$(B_K)_{E,E'} = \int_K w_{K,E}^T \mathcal{K}_K^{-1} w_{K,E'} dx.$$

▶ The **mixed hybrid** formulation necessitates inverting  $B_K$ .

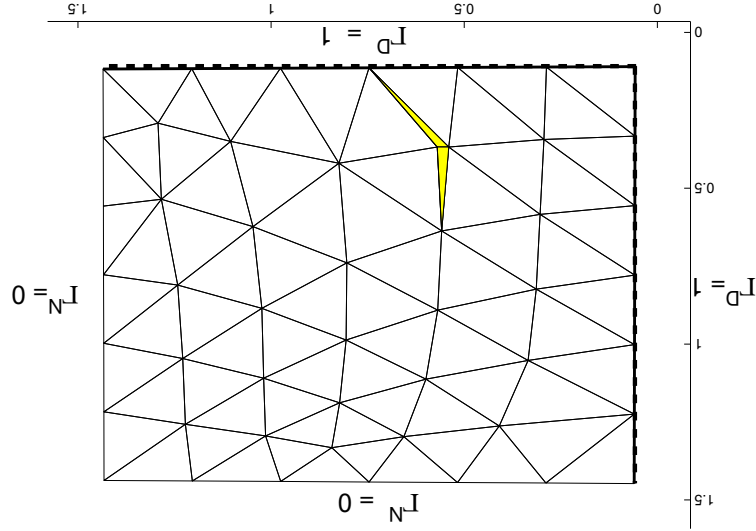
▶ The stability of the numerical code is crucial due to possible ill-conditioned matrices

▶ Two codes are used to invert  $B_K$

1. Cramer's method (generated automatically by **Maple** )
2.  $LDL^T$  Factorization (  $B_K^{-1} = L^{-T} D^{-1} L^{-1}$  )

## Effect of the spatial discretization

Numerical test:



$$T = 3, \Delta t = 1, s = 1, \kappa = I_{2 \times 2}, f = 0.$$

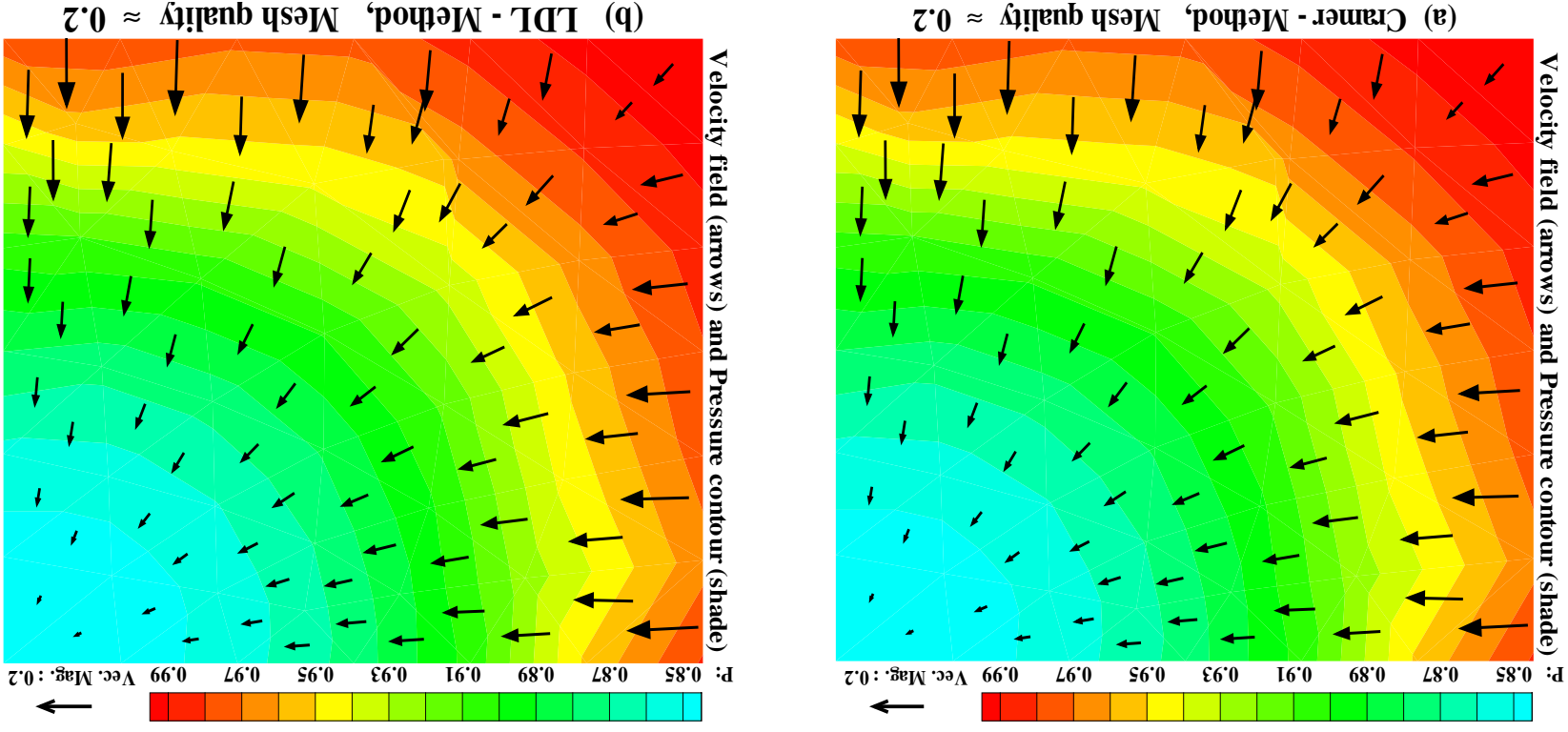
The quality of a triangular element  $K$  is evaluated by

$$\mathcal{Q}(K) = 2\sqrt{3} \frac{h_K}{\rho_K}, \quad \mathcal{Q}(K) = 1 \text{ if } K \text{ equilateral triangle}$$

$\rho^K$ : is the element diameter (the length of its longest edge).  
 $h^K$ : is the in-radius.

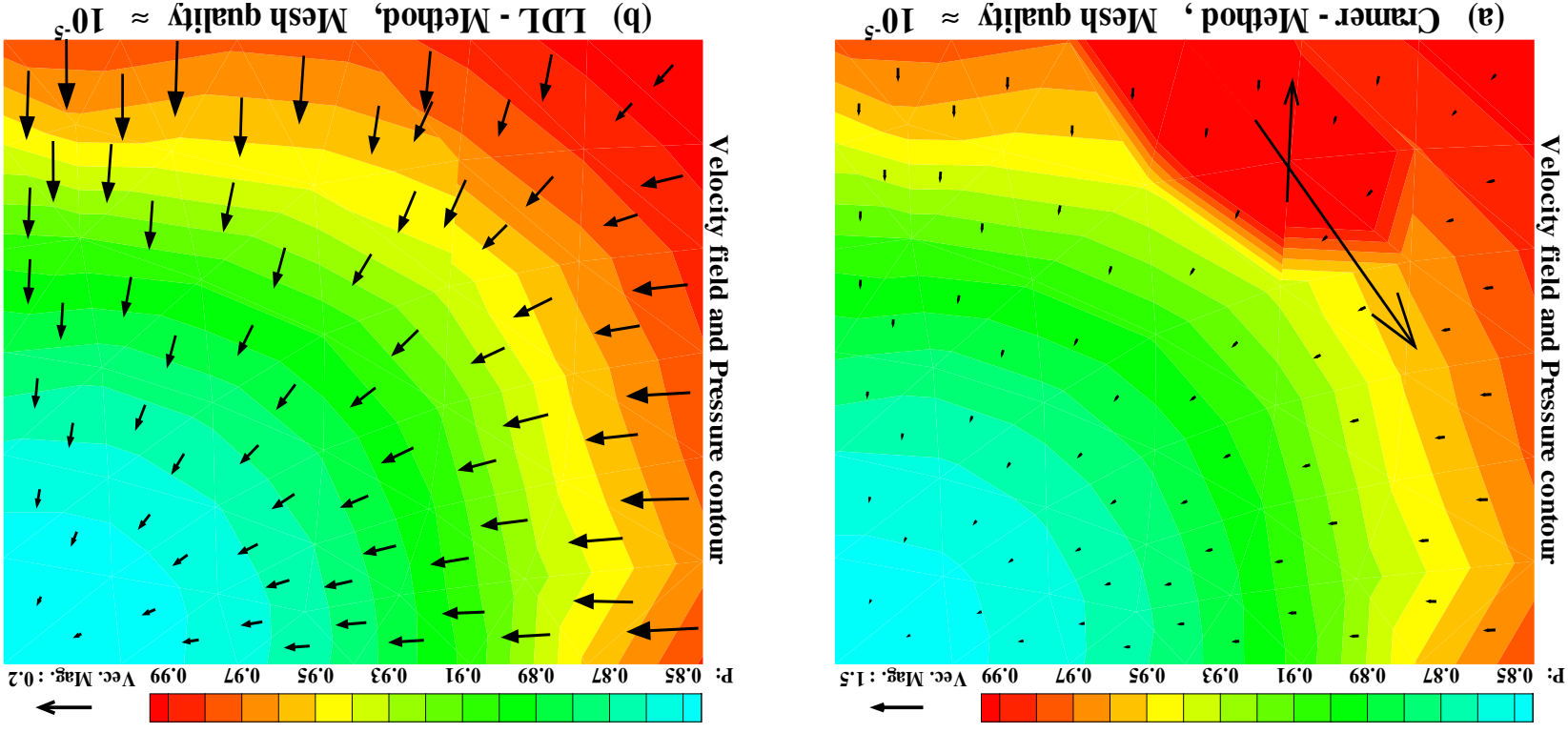
Comparison between 2 computations obtained by using the 2 codes  
in the MHFE algorithm

well-shaped,  $\mathcal{Q}(K) \approx 0.2$



## Velocity field and pressure contour

Comparison between 2 computations obtained by using the 2 codes  
in the MFE algorithm  
ill-shaped,  $\mathcal{Q}(K) \approx 10^{-5}$

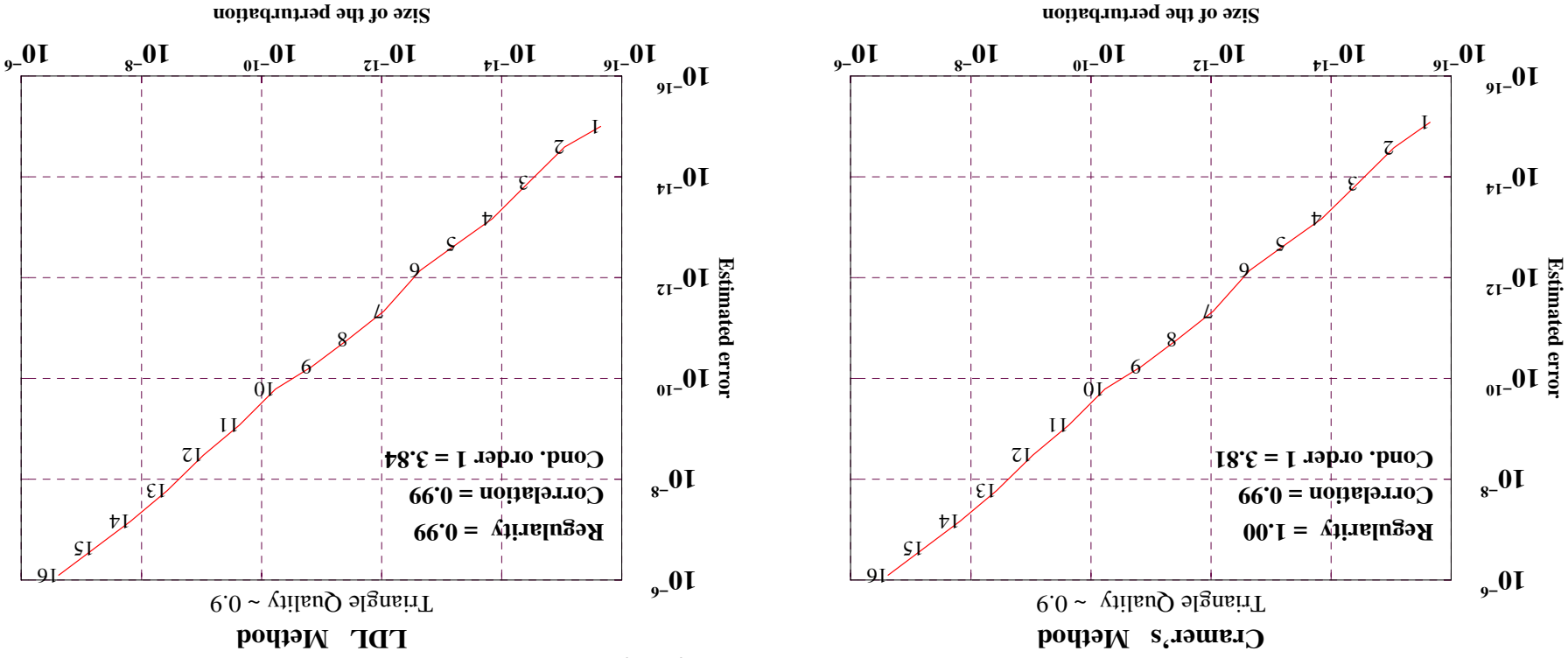


Velocity field and pressure contour

# Effect of the spatial discretization

Stability analysis by Aquarels

well-shaped,  $\mathcal{Q}(K) \approx 0.9$



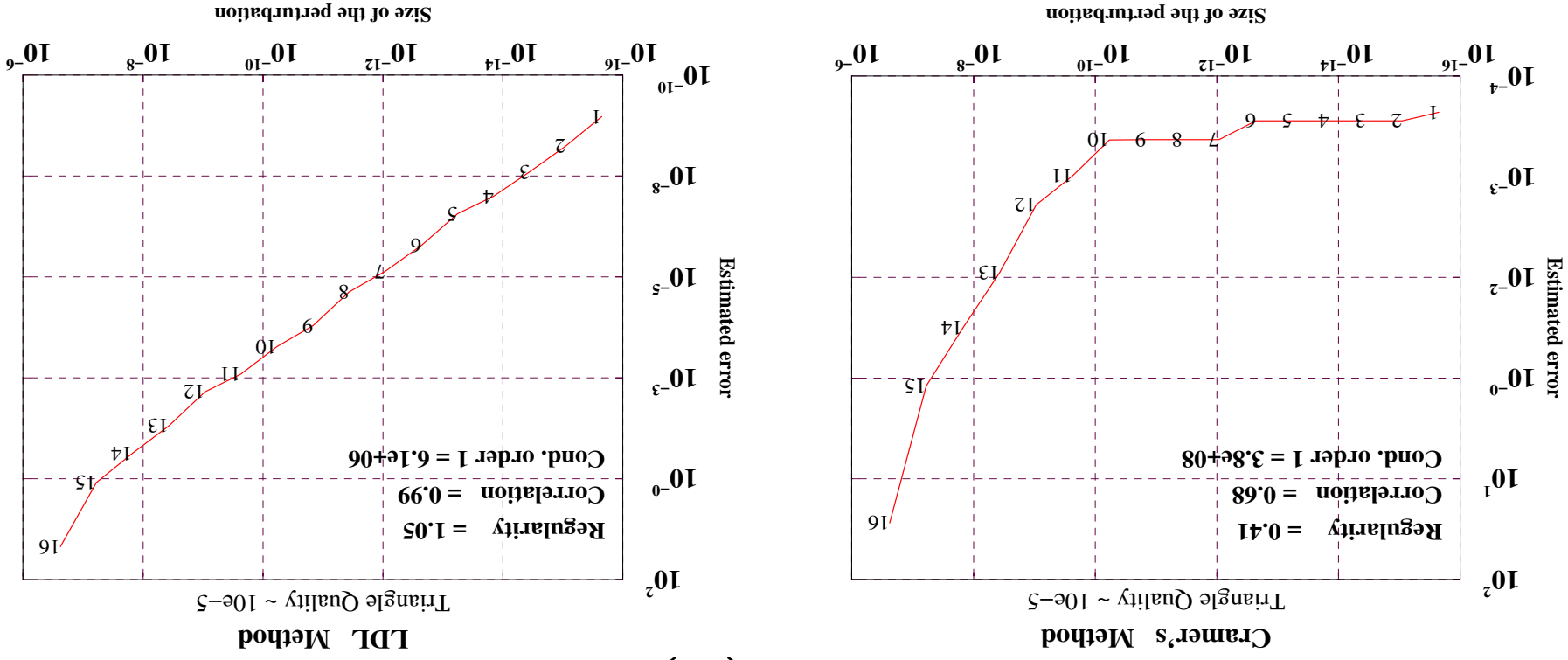
Absolute error $\ B_K^{-1} - \tilde{B}_K^{-1}\ _2$	Cramer	$1.7 \times 10^{-15}$
	LDL	$4.4 \times 10^{-16}$

Relative error	Cramer	$5.9 \times 10^{-16}$
	LDL	$1.5 \times 10^{-16}$

# Effect of the spatial discretization

## Stability analysis by Aquarels

||-shaped,  $\mathcal{Q}(K) \approx 10^{-5}$

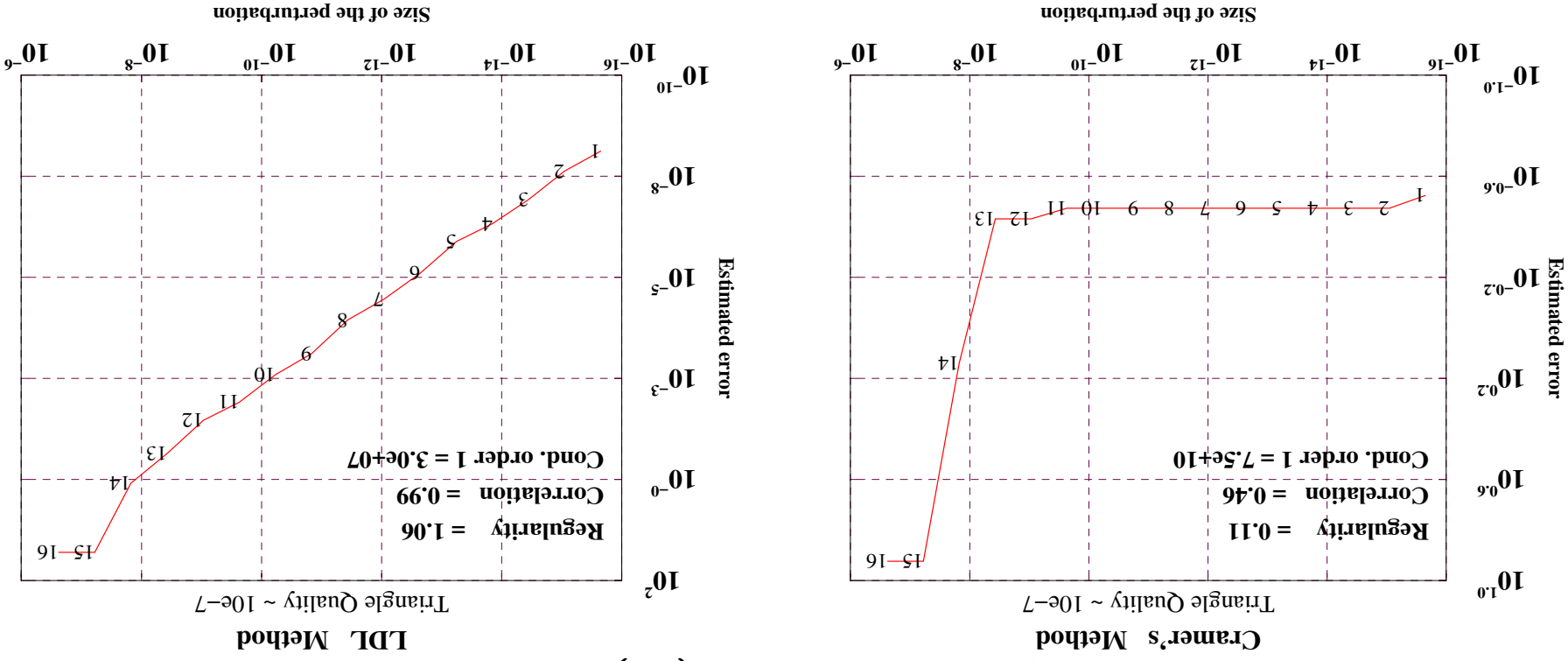


Method	Absolute error	Relative error
Cramer	$1.6 \times 10^{+03}$	$1.1 \times 10^{-04}$
LDL	$4.7 \times 10^{-03}$	$3.3 \times 10^{-10}$

# Effect of the spatial discretization

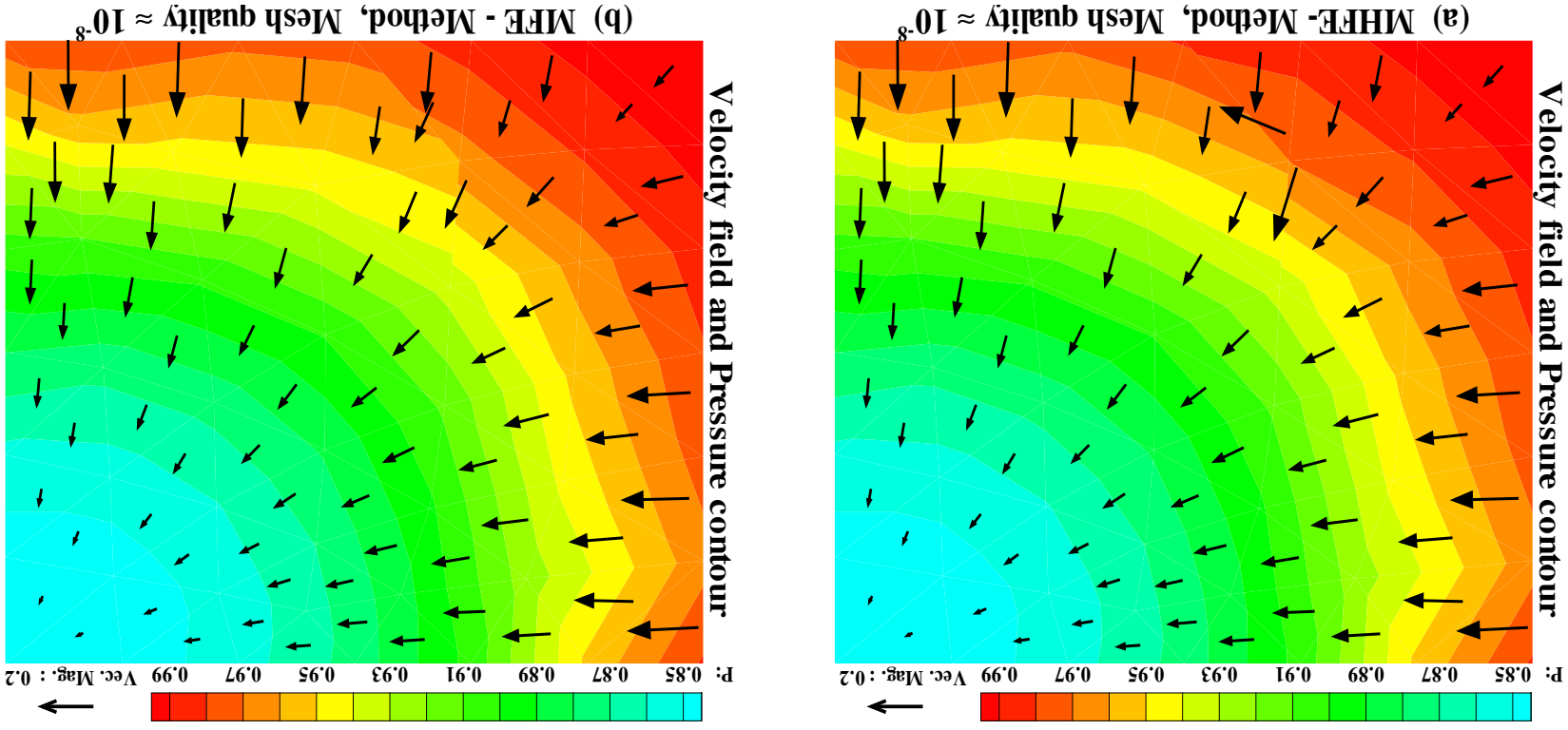
Stability analysis by Aquarels

very ill-shaped,  $\mathcal{Q}(K) \approx 10^{-7}$



Method	Relative error	Absolute error
LDL	$1.2 \times 10^{-10}$	$3.5 \times 10^{+07}$
Cramer	$1.6 \times 10^{-01}$	$2.7 \times 10^{-02}$

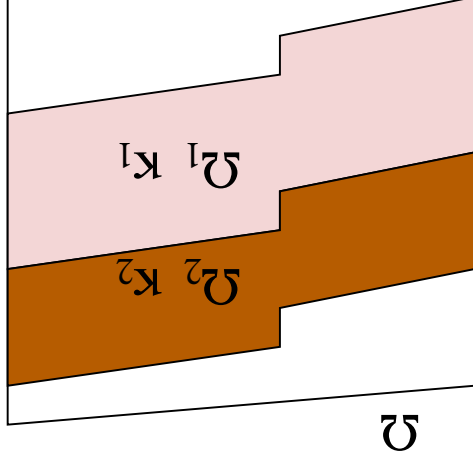
Comparison between MHFE et MFE methods  
 $\mathcal{Q}(K) \approx 10^{-8}$ , Condition number  $\approx 10^{15}$



Velocity field and pressure contour

# Behavior of MFE and MHFE methods on heterogeneous media

Suppose that the heterogeneous domain  $\Omega$  is composed of two sub-domains  $\Omega_i$ , according to their conductivities  $K_i = \kappa_i I_{2 \times 2}$



Decomposition of the domain according to the permeability values.

with very high contrast of permeability:  $\frac{\kappa_2}{\kappa_1} \gg 1$

**Lemma 1.**

Let  $A$  (resp.  $B$ ) be a **non-singular** (resp. **singular**) matrix, then the condition number  $\kappa(A)$  of  $A$  is bounded by the following inequality

$$\kappa(A) = \|A\| \cdot \|A^{-1}\| \geq \frac{\|A\|}{\|A - B\|}$$

**Proof:**

Since  $B$  is a singular matrix, there exists a nonzero vector  $x$  such that  $Bx = 0$ , then we can write

$$\begin{aligned} \|Ax\| &= \|(A - B)x\| \leq \|A - B\| \|x\|, \\ \|x\| &= \|A^{-1}Ax\| \leq \|A^{-1}\| \|Ax\|. \end{aligned}$$

Hence, the lemma is a direct consequence of these inequalities.

Let  $\mathcal{J}_{k_1, k_2}$  (resp.  $\mathcal{J}_{k_1}$ ) denote the matrix  $\mathcal{J}$  over  $\Omega = \Omega_1 \cup \Omega_2$  ( resp.  $\Omega = \Omega_1$ )

$$\mathcal{J} = \begin{pmatrix} -R \\ D & -R^T \\ M \end{pmatrix}$$

**Proposition 1.**  $\kappa(\mathcal{J}_{k_1, k_2}) = \mathcal{O}(\frac{k_1}{k_2})$ .

**Proof:** We have

$$\mathcal{J}_{k_1, k_2} = \begin{cases} \mathcal{J}_{k_1} & \text{in } \Omega_1, \\ \mathcal{J}_{k_2} & \text{in } \Omega_2, \\ \mathcal{J}_{k_1, k_2} & \text{on } \Omega_1 \cap \Omega_2. \end{cases}, \text{ define } \underline{\mathcal{J}}_{k_1, k_2} = \begin{cases} 0 & \text{in } \Omega_1, \\ \mathcal{J}_{k_2} - \mathcal{J}_{k_1} & \text{in } \Omega_2, \\ \mathcal{J}_{k_1, k_2} - \mathcal{J}_{k_1} & \text{on } \Omega_1 \cap \Omega_2. \end{cases}$$

$\underline{\mathcal{J}}_{k_1, k_2}$  is singular with  $\mathcal{J}_{k_1, k_2} - \underline{\mathcal{J}}_{k_1, k_2} = \mathcal{J}_{k_1}$

Hence  $\kappa(\mathcal{J}_{k_1, k_2}) \geq \frac{\|\mathcal{J}_{k_1, k_2}\|}{\|\mathcal{J}_{k_1}\|} = \mathcal{O}(\frac{k_1}{k_2})$ ,

By applying **Lemma.1**, the following estimations of the condition numbers can be shown:

**MFE** Schur complement matrix (unknowns  $Tp$ )

$$1. \kappa(A_{\kappa_1, \kappa_2}) = \mathcal{O}\left(\frac{\kappa_1}{\kappa_2}\right)$$

**MFE** Schur complement matrix (unknowns  $Q$ )

$$2. \kappa(\tilde{A}_{\kappa_1, \kappa_2}) = \mathcal{O}\left(\max\left\{\kappa_1 \Delta t \|S^{-1}\|, \frac{\kappa_1}{\kappa_2} \left(\frac{1 + \kappa_1 \Delta t \|S^{-1}\|}{1 + \kappa_2 \Delta t \|S^{-1}\|}\right)\right\}\right)$$

**MFE** matrix without eliminating the fluxes (unknowns  $Q$  and  $P$ )

$$3. \kappa(\mathcal{L} - \Delta t \tilde{f}_{\kappa_1, \kappa_2}) = \mathcal{O}\left(\max\left\{\frac{\kappa_2}{\kappa_1} \|\Delta t\|, \frac{\kappa_1}{\kappa_2}\right\}\right)$$

## Accumulation of numerical errors by the MHFEM

After computing the pressure and its traces, the fluxes are computed locally by solving a block diagonal system

$$B \hat{Q} = (P - T^P). \quad (11)$$

$B$  is a block diagonal matrix of condition number  $\kappa(B) = \mathcal{O}(\frac{\kappa_1}{\kappa_2})$ . Thus, the fluxes  $\hat{Q}$  are computed with conditioning order  $\mathcal{O}(\left(\frac{\kappa_1}{\kappa_2}\right)^2)$ .

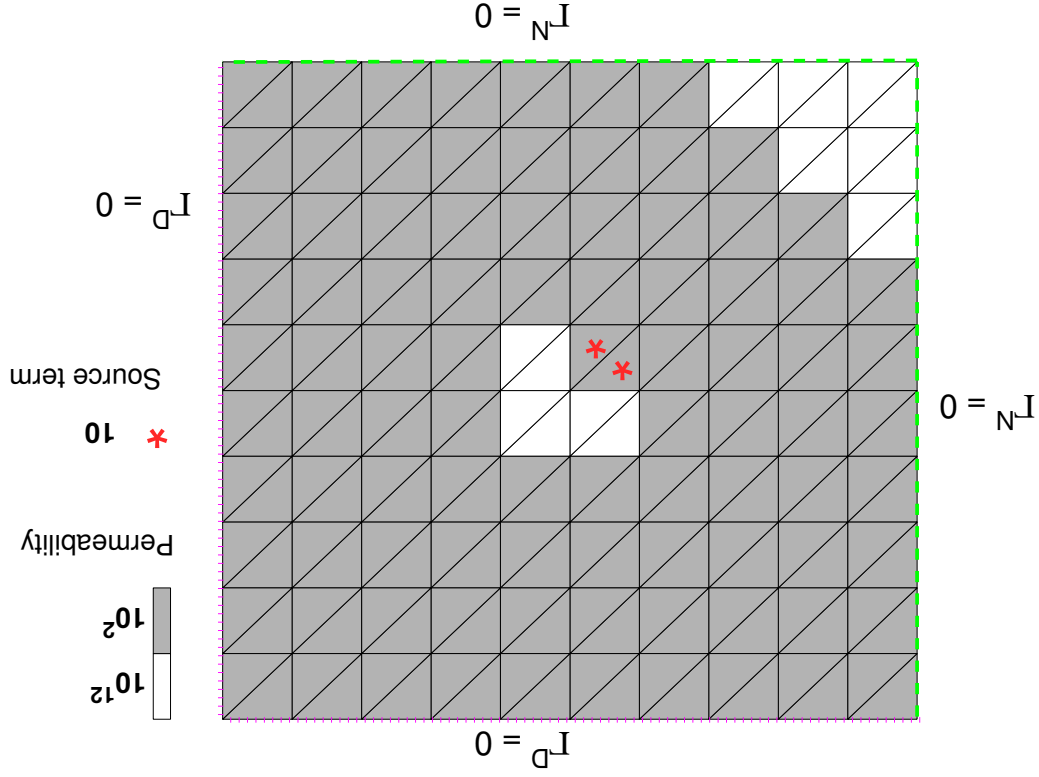
In the stationary case, the absolute roundoff error of the local mass balance can be written as

$$\left[ \sum_{E \subset \partial K} q_{K,E} - f_K \right] = \kappa_K \sum_{E \subset \partial K} \epsilon_{K,E}. \quad (12)$$

where  $\epsilon_{K,E}$  denotes the numerical error produced from the computation of  $P$  and  $T^P$ .

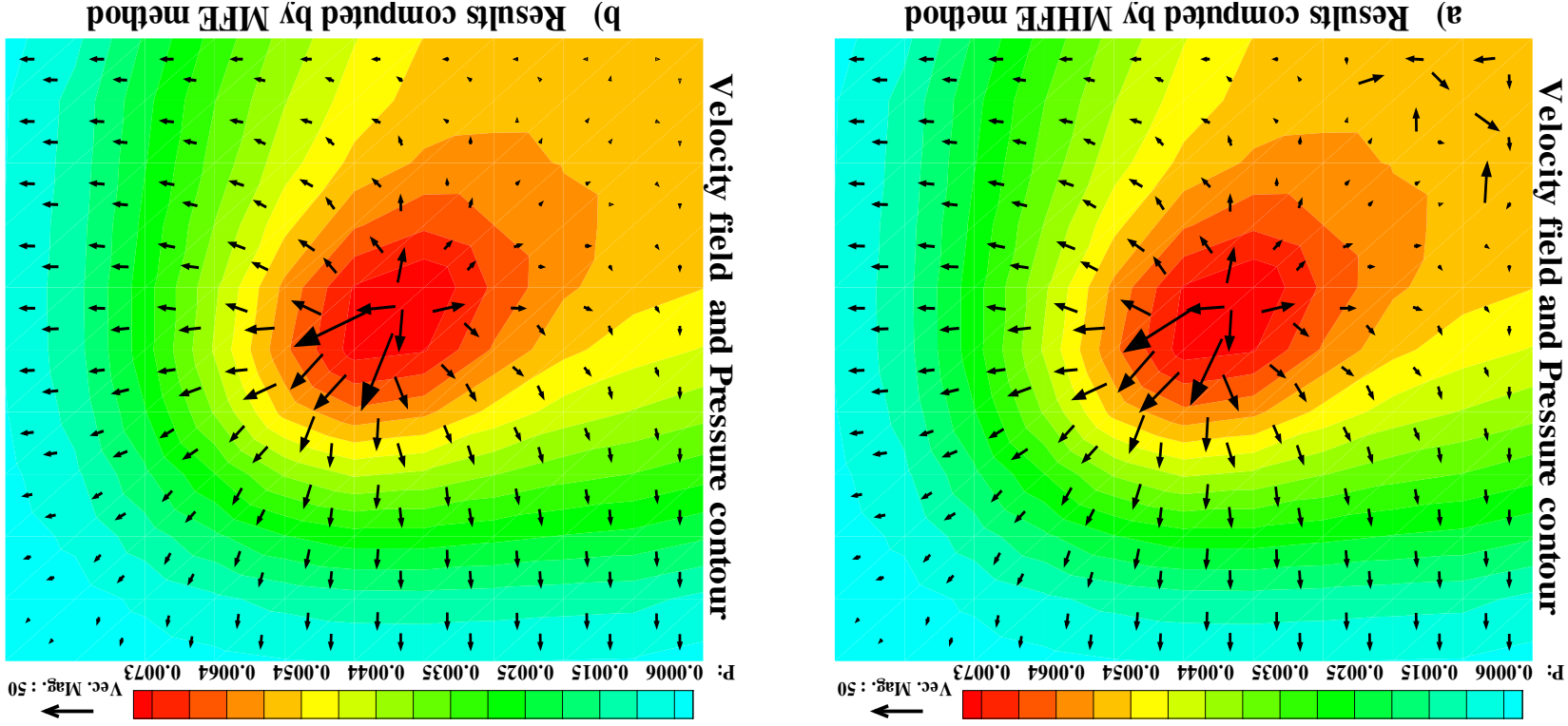
Experiment 1: stationary problem

Consider a stationary problem over a unit square with the following initial and boundary conditions.



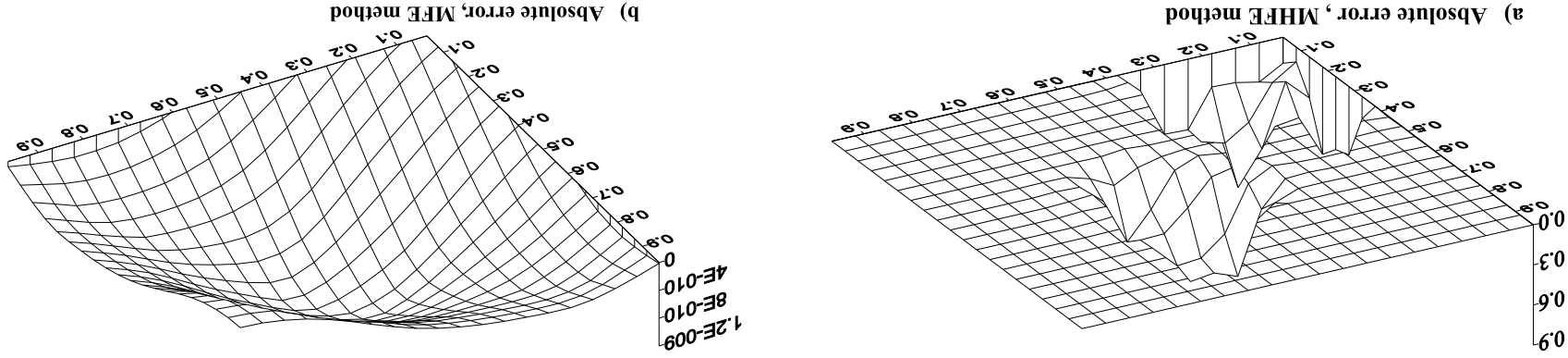
Triangulation of the unit square.

Experiment 1: stationary problem

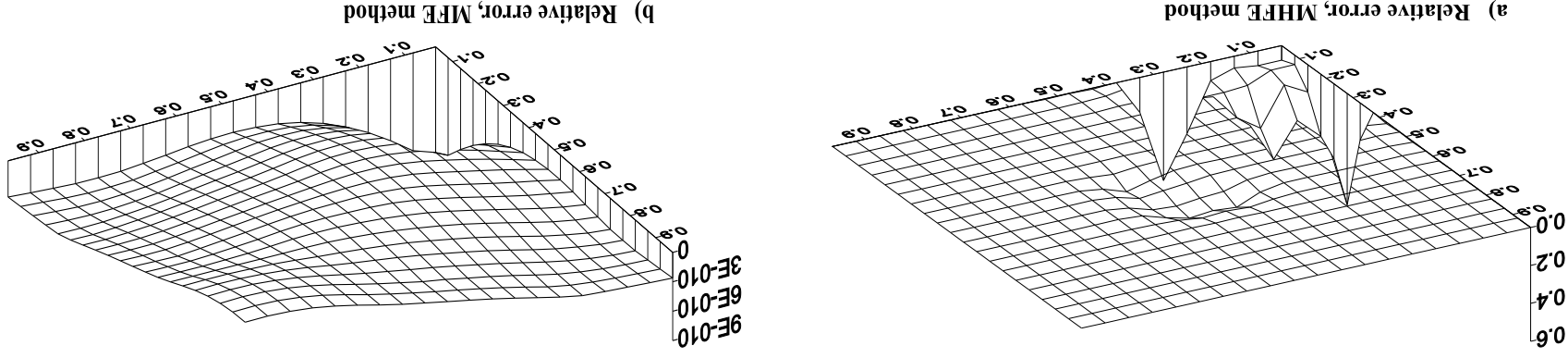


Results obtained by the MHE and MFE methods

Experiment 1: stationary problem



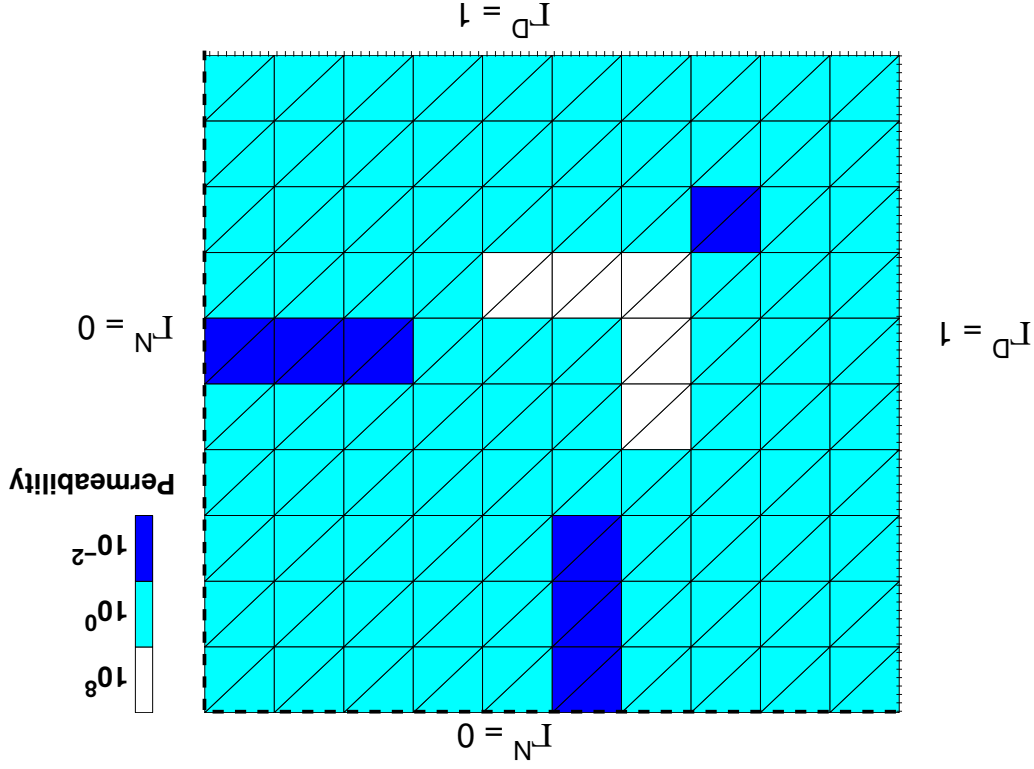
Absolute errors of the mass balance over the domain



Relative errors of the mass balance over the domain

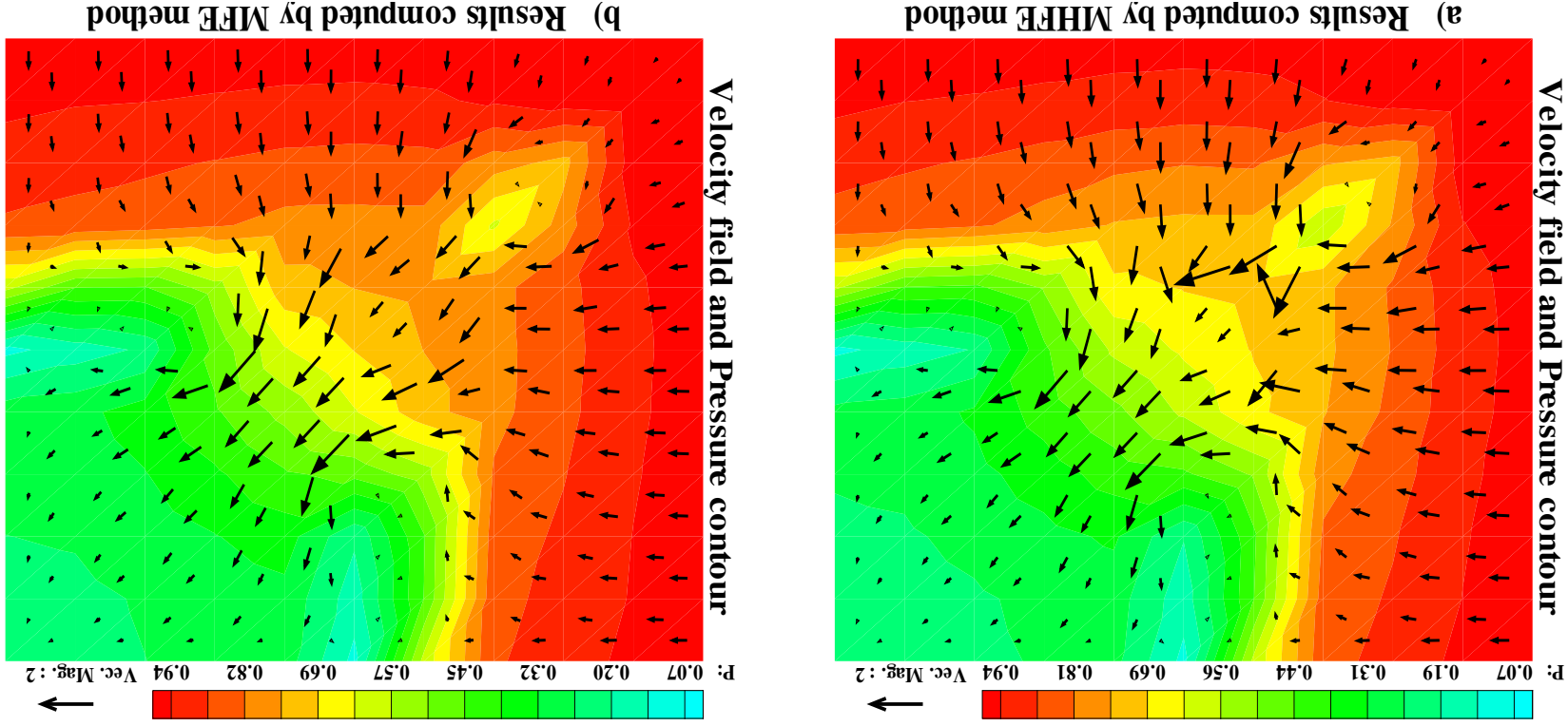
Experiment 2: transient problem

Consider a time-dependent problem over a unit square with the following initial and boundary conditions.  $\mathcal{T} = [0, 1]$ ,  $\Delta t = \frac{1}{10}$ ,  $s = 1$ ,  $f = 0$ .



Permeability distribution and boundary conditions

Experiment 2: transient problem



Results obtained by the MHFE and MFE methods

**Eisenstat** preconditioning is used with the CG algorithm to solve the S.P.D systems.

The symmetric indefinite systems are solved by using **Symmlq** solver without any preconditioning.

## Computing times

$\Omega = \Omega_1 \cup \Omega_2, \Delta t = 1/10, s_i = 1 \text{ } i = 1, 2, \text{ PCG solver.}$

$\kappa_2/\kappa_1$	MHFE method			MFE method		
	CPU Cond.	Num.	# Iter.	CPU Cond.	Num.	# Iter
$10^2$	2.2	$10.2 \times 10^2$	22-21	1.8	$4.3 \times 10^2$	33-31
$10^4$	2.3	$10.9 \times 10^4$	24-23	2.1	$4.2 \times 10^4$	42-40
$10^6$	2.4	$10.7 \times 10^6$	26-25	2.4	$4.5 \times 10^6$	55-54

$\Delta t = 1/10, \kappa_i = 1 \text{ } i = 1, 2, \text{ PCG solver.}$

s	MHFE method			MFE method		
	CPU Cond.	Num.	# Iter.	CPU Cond.	Num.	# Iter.
$10^2$	2.2	50.9	20-19	2.4	$65.1 \times 10^1$	30-29
$10^{-4}$	2.2	55.1	20-19	3.2	$68.3 \times 10^3$	51-50
$10^{-6}$	2.2	55.1	20-19	3.9	$68.3 \times 10^5$	84-83

- ▷ The **MHFE** and **MFE** methods lead to symmetric, positive definite linear systems for parabolic problems.
- ▷ In heterogeneous media, the condition numbers of both systems grow up according to the ratio between the highest and the lowest values of permeability parameters between adjacent subdomains.
- ▷ The **MHFE** formulation sometimes leads to senseless values in the velocity field, especially when large jumps in the hydraulic conductivity take place.
- ▷ Small values of the storage coefficient could blow up the conditioning of the **MFE** linear system. Thus, the symmetric indefinite linear system has to be solved.