# Hermitian matrices : Spectral coupling, plane geometry/trigonometry and optimisation 

Françoise Chatelin and<br>M. Monserrat Rincon-Camacho

Technical Report TR/PA/15/31

Publications of the Parallel Algorithms Team
http://www.cerfacs.fr/algor/publications/

## £CERFACS

## ( Qualitative

# HERMITIAN MATRICES : SPECTRAL COUPLING, PLANE GEOMETRY/TRIGONOMETRY AND OPTIMISATION 

FRANÇOISE CHATELIN ${ }^{(1),(2)}$ AND M. MONSERRAT RINCON-CAMACHO ${ }^{(1)}$


#### Abstract

The paper presents the information processing that can be performed by a general hermitian matrix when two of its distinct eigenvalues are coupled, such as $\lambda<\lambda^{\prime}$. Setting $a=\frac{\lambda+\lambda^{\prime}}{2}$ and $e=\frac{\lambda^{\prime}-\lambda}{2}>0$, the new spectral information that is provided by coupling is expressed in terms of the ratios $\frac{e}{|a|}\left(\right.$ if $\left.\lambda \lambda^{\prime}>0\right)$ or $\frac{|a|}{e}\left(\right.$ if $\left.\lambda \lambda^{\prime}<0\right)$ and of the product $|a| e$. The information is delivered in geometric form, both metric and trigonometric, associated with various right-angled triangles deriving from optimality conditions. The paper contains a generalisation to indefinite matrices over $\mathbb{R}$ or $\mathbb{C}$ of Gustafson's operator trigonometry which in the matrix case assumes definiteness (mostly over $\mathbb{R}$ ).


Keywords: Spectral coupling, hermitian matrix, indefinite symmetric or hermitian, spectral plane, observation angle, perspective angle, catchvector, antieigenvector, middle vector, Euler equation, balance equation, torus in 3D.

## 1. Composite natural phenomena

1.1. Introduction. It is well-known that many natural phenomena are composite: they are produced by the interaction of two (or more) simpler phenomena, where the interaction is in general nonlinear. The coupling of two phenomena is ubiquitous in Physics; electromagnetism, convection-diffusion, plasma fusion, flutter are just a few examples. And it is known that the tighter the coupling, the higher the departure of normality for the matrix resulting from discretisation [Chaitin-Chatelin and Frayssé, 1996, chapter 10]. The theoretical understanding and algorithmic resolution of tightly coupled phenomena remains today a formidable challenge for scientists. A classical path of investigation is to consider the linear coupling $A+t B=A(t)$ by the complex parameter $t$ between two independent phenomena modelled by $A$ and $B$. Physical examples of such a parameter are given by Reynolds and Péclet numbers, which are real or by the admittance which can be complex. The evolution under $t$ in $\mathbb{C} \cup\{\infty\}$ of the spectrum of $A(t) \in \mathbb{C}^{n \times n}$ is studied in [Chatelin, 2012b, chapter 7]. This approach, known as Homotopic Deviation, puts no restriction on $\|t B\|=|t|\|B\|$, which, by contrast, is classically assumed to be small enough in perturbation theory.

In this paper, we propose a different angle of approach to coupling based on two real parameters related to a single matrix $A$ enjoying a property of symmetry that we describe next.
1.2. Spectral coupling. In the work we present below, we focus our attention on the coupling of any two distinct real eigenvalues $\lambda<\lambda^{\prime}$ of a general hermitian or symmetric matrix $A$, a coupling called spectral coupling. Such a coupling can be seen as a self-interference for $A$, that is an interference of $A$ with itself by means of two of its eigenvalues. Hence we develop new consequences of the classical spectral theory for symmetric matrices [Parlett, 1998] and hermitian matrices [Horn and Johnson, 1985, chapter 4, Chatelin, 2012a, chapter 8] These consequences are almost all derived from the fundamental property that the quadratic form $x^{H} A x$ is real for $x \in \mathbb{C}^{n}$. Of special interest are the new informations provided by the simultaneous consideration of two eigenelements (values/vectors) rather than a single one in classical theory.

[^0]1.3. Related work. The systematic study of spectral coupling that we have developped originates in the pioneering work that Gustafson has pursued in relative isolation for almost five decades [Gustafson, 2012]. His context deals more generally with strongly accretive operators in Hilbert spaces, which admit symmetric or hermitian positive definite matrices as particular cases. The work we present diverges and deals more generally with indefinite matrices.

When referring to his own theory, Gustafson uses several terms such as operator trigonometry, antieigenvalue analysis, introducing the concepts of anti-eigen(-value/-vector). The reason why his original vocabulary has to be adapted to the context of indefinite matrices will become clear later on. Also, we shall indicate in what sense Gustafson's "antieigenvectors" can be interpreted as the "most turned" vectors by $A$ when the ground field $K$ is $\mathbb{C}$ rather that $\mathbb{R}$.
1.4. Organisation of the paper. Section 2 presents the new information provided, in the numerical plane $\mathbb{R}^{2}$, by the simultaneous consideration of 2 real eigenvalues $\lambda<\lambda^{\prime}$ for a matrix $A$ which is either symmetric $\left(A=A^{T} \in \mathbb{R}^{n \times n}\right)$ or hermitian $\left(A=A^{H} \in \mathbb{C}^{n \times n}\right)$. The two isolated eigenvalues are linked by the cercle $\Gamma$ centered at the mean $\left(\frac{\lambda+\lambda^{\prime}}{2}, 0\right)$ with radius $\frac{\lambda^{\prime}-\lambda}{2}$ drawn in $\mathbb{R}^{2}$. The proofs involve only elementary plane geometry and trigonometry in a triangle. The consideration of two associated orthonormal eigenvectors $q$ and $q^{\prime}$ in $K^{n}$ may require to distinguish whether the ground field $K$ is $\mathbb{R}$ or $\mathbb{C}$. Section 3 is devoted to the case $K=\mathbb{R}$ and to trigonometry in the invariant real plane $\mathbb{R}^{2}$ spanned by $q$ and $q^{\prime}$ in $\mathbb{R}^{n}$. Three variational principles are derived in Section 4 by looking at the stationary values of a more generic functional $c: K^{n} \rightarrow \mathbb{R}$ (due to the second author). Next Section 5 presents a geometric interpretation in $\mathbb{R}^{3}$ of the optimality principles obtained in the space $\mathbb{C}^{2} \cong \mathbb{R}^{4}$ (when $K=\mathbb{C}$ ) with 2 complex dimensions, i.e. 4 real ones. Section 6 illustrates some of the underrated properties of the "middle vectors" that are the linear combinations $z q+z^{\prime} q^{\prime}$ with $|z|=\left|z^{\prime}\right|=\frac{1}{\sqrt{2}}, z, z^{\prime} \in \mathbb{C}$. Two examples from Numerical Analysis and Statistics are provided. Section 7 deals with a second generic functional $\nu: \mathbb{R} \times K^{n} \rightarrow \mathbb{R}^{+}$. The paper closes on Section 8 which wraps up the information processing provided in $\mathbb{R}^{2}$ by spectral coupling, that is explained by the underlying algebraic structure of $\mathbb{R}^{2}$ as the ring of bireal numbers.

## 2. Spectral information processing in a spectral plane

2.1. Definition. Let $A \in K^{n \times n}, K=\mathbb{R}$ or $\mathbb{C}$, be a symmetric ( $K=\mathbb{R}, A=A^{T}$ ) or hermitian ( $K=\mathbb{C}, A=A^{H}$ ) matrix. The spectrum of $A$ consists of $n$ real eigenvalues $\lambda_{1} \leq \ldots \leq \lambda_{n}$ lying on the spectral line $\mathbb{R}$. If $A \neq \lambda I$, the spectrum contains at least two distinct eigenvalues. The matrix $A$ is diagonalisable in the eigenbasis $Q=\left[q_{1}, \ldots, q_{n}\right]$, with $Q^{-1}=Q^{H}$ if $K=\mathbb{C}$ and $Q^{-1}=Q^{T} \in \mathbb{R}^{n \times n}$ if $K=\mathbb{R}$.

Let us turn to the simultaneous consideration of two distinct eigenvalues $\lambda<\lambda^{\prime}$ which produces new information (both geometric and trigonometric) best understood if one considers the spectral plane $\mathbb{R}^{2}$ denoted $\boldsymbol{\Sigma}$ which contains the spectral line $\mathbb{R}$ as its real axis.
2.2. Notation for the eigenvalue pair $\left\{\lambda, \lambda^{\prime}\right\}$. The eigenvalues $\lambda$ and $\lambda^{\prime}$ are the distinct real roots of the quadratic equation

$$
\begin{equation*}
\mu^{2}-2 a \mu+g^{2}=0, \quad a=\frac{\lambda+\lambda^{\prime}}{2}, \quad g^{2}=\lambda \lambda^{\prime} \tag{2.1}
\end{equation*}
$$

where $a$ is the arithmetic mean. By assumption $a^{2}-g^{2}=\frac{1}{4}\left(\lambda-\lambda^{\prime}\right)^{2}>0$. We set $e=\frac{\lambda^{\prime}-\lambda}{2}>0$ and observe that $a^{2}=g^{2}+e^{2} \Leftrightarrow-e^{2} \leq g^{2}<a^{2}$. Moreover $g^{2}=0 \Leftrightarrow \lambda$ or $\lambda^{\prime}=0, \lambda^{\prime}$ or $\lambda=2 a$ $\Leftrightarrow e=|a|$.
2.3. The three pythagorean means of $\lambda$ and $\lambda^{\prime}, 0<\lambda<\lambda^{\prime}$. To the two positive numbers $0<\lambda<\lambda^{\prime}$ are associated three types of mean: arithmetic $a=\frac{\lambda+\lambda^{\prime}}{2}$, geometric $g=\sqrt{\lambda \lambda^{\prime}}$, harmonic $h=\frac{g^{2}}{a}=\frac{2 \lambda \lambda^{\prime}}{\lambda+\lambda^{\prime}}$. The three means which satisfy $0<\lambda<h<g<a<\lambda^{\prime}$ are known since Antiquity as the pythagorean means.

We recall that $g^{2}=\lambda \lambda^{\prime}$ is always positive, i.e. $|a|>e$, when $A$ is definite, leading to $g=\sqrt{\lambda \lambda^{\prime}}$. By contrast, when $A$ is indefinite, $\lambda \lambda^{\prime}$ may be nonpositive $(|a| \leq e)$ leading to $g=0$ $(|a|=e)$ or $|g|=\sqrt{-\lambda \lambda^{\prime}}>0(|a|<e)$.
2.4. The triangle $\operatorname{Tr}(M)=O M C$ in the spectral plane $\boldsymbol{\Sigma}$. Let be given the pair $\left\{\lambda, \lambda^{\prime}\right\}$, $\lambda<\lambda^{\prime}$ lying on the spectral line. We consider the circle $\Gamma$ centered at $C, \overline{O C}=a$ with radius $e$, which passes through the points $\Lambda=(\lambda, 0)$ and $\Lambda^{\prime}=\left(\lambda^{\prime}, 0\right)$ and lies in the spectral plane. Depending on the sign of $g^{2}=\lambda \lambda^{\prime}$, the origin $O$ is outside $\Gamma\left(g^{2}>0\right)$, on $\Gamma(g=0)$ or inside $\Gamma$ $\left(g^{2}<0\right): g^{2}$ is the power of $O$ with respect to $\Gamma$. The circle $\Gamma$ can be thought of as a linking curve between the isolated eigenvalues $\lambda$ and $\lambda^{\prime}$, a curve which necessarily lies in a plane where elementary geometric constructions and trigonometric calculations can be performed.

We first suppose that $g^{2} \neq 0$, see Figure 1.


Figure 1. $\operatorname{Tr}(M)$ when $g^{2} \neq 0$ in $\boldsymbol{\Sigma}$

When $O$ is inside (resp. outside) $\Gamma$, an arbitrary line drawn through $O$ intersects $\Gamma$ in two points $M$ and $N$ always (resp. if the acute angle $\alpha=\angle(O C, O M)$ is not too large). We consider the triangle $\operatorname{Tr}(M)=O M C$ defined for any $M$ on $\Gamma, M \neq \Lambda$ or $\Lambda^{\prime}$, and $a \neq 0(O \neq C)$.

Two of the side lengths are fixed: $O C=|a|$ and $M C=e$, while the third length $O M$ varies with $M$. We denote the three ordinary angles of $\operatorname{Tr}(M)$ as follows: $\alpha=\angle(O C, O M)$, $\beta=\angle(M C, M O)$ and $\gamma=\angle(C O, C M)$ where the angles vary in $] 0, \pi\left[\right.$ according to $\operatorname{sgn}\left(g^{2}\right)$ as follows:

| $g^{2}$ | - | + |
| :---: | :---: | :---: | :---: |
| $\alpha$ | $\alpha \in] 0, \pi[$ | $0<\alpha<\frac{\pi}{2}$ |
| $\beta$ | $0<\beta<\frac{\pi}{2}$ | $\beta \in] 0, \pi[$ |
| $\gamma$ | $0<\gamma<\pi$ |  |

For future reference, we also introduce $\delta=\frac{\gamma}{2}=\angle\left(\Lambda^{\prime} \Lambda, \Lambda^{\prime} M\right), 0<\delta<\frac{\pi}{2}$. See Figure 1. We recall that $\alpha+\beta+\gamma=\pi$ and $\frac{\sin \alpha}{e}=\frac{\sin \beta}{|a|}=\frac{\sin \gamma}{O M}$, hence the ratio $\frac{\sin \alpha}{\sin \beta}=\frac{e}{|a|}$ is fixed.


Figure 2. Trigonometry for $\alpha$ (a) and $\beta$ (b) in $\boldsymbol{\Sigma}$

If $a=0, O=C$ in case (b) and the triangle is degenerate $(\beta=0)$.
Lemma 2.1. (a) When $0<g^{2}<a^{2}, 0<\alpha \leq \phi$ with $\sin \phi=\frac{e}{|a|}$, and $\alpha=\phi<\frac{\pi}{2} \Leftrightarrow \beta=\frac{\pi}{2}$, $\gamma=\frac{\pi}{2}-\phi$.
(b) When $-e^{2}<g^{2}<0, \beta \leq \psi$ with $\sin \psi=\frac{|a|}{e}$, and $\beta=\psi \Leftrightarrow \alpha=\frac{\pi}{2}, \gamma=\frac{\pi}{2}-\psi$.

Proof. Elementary trigonometry. (a) $C H \leq C T=e$. The angle $\alpha$ is maximum at the value $\phi$ when the secant line $O M$ is tangent to $\Gamma$ at $T$ or $U$. Then $O T^{2}+e^{2}=a^{2}, O T=O U=g$, $\cos \phi=\frac{o T}{|a|}=\frac{g}{|a|}=\sin \gamma$.
(b) $C H \leq C O=|a|>0$. The angle $\beta$ is maximum at the value $\psi$ when the line $O M$ is orthogonal to the spectral axis, and intersects $\Gamma$ at $V$ and $W$. Then $O V^{2}+a^{2}=e^{2}$ and $O V=O W=|g|, \cos \psi=\frac{|g|}{e}=\sin \gamma$.

When $g^{2}>0$, the circle $\Gamma$ is seen from the origin $O$ under the angle $2 \phi$, hence $\phi$ is the (outer) observation angle. When $g^{2}<0, a \neq 0$, the segment $O C \neq 0$ is seen from the circle $\Gamma$ under the angle $\beta$ : the maximal angle $\psi$ is called the (inner) perspective angle. If $g^{2} \rightarrow 0^{+}$(resp. $\left.0^{-}\right) \phi$ (resp. $\psi$ ) tends to $\frac{\pi}{2}$, hence $\gamma \rightarrow 0$. The points $T, U$ (resp. $V, W$ ) tend to $O$ and the right-angled triangles tend to collapse into the segment $O C$.

We observe that $\cos \phi=\frac{g}{|a|}$ or $\cos \psi=\frac{|g|}{e}$ for $M$ at two locations on $\Gamma$ : either $T$ and $U$ $\left(\beta=\frac{\pi}{2}, O T=g, g^{2}>0\right)$ or $V$ and $W\left(\alpha=\frac{\pi}{2}, O V=|g|, g^{2}<0\right)$.

By contrast the value $\frac{e}{|a|}$ (resp. $\frac{|a|}{e}$ ) is fixed, independently of the location of $M$ on $\Gamma$. Of course the number can be interpreted as the maximum value $\sin \phi$ (resp. sin $\psi$ ) only if $\operatorname{Tr}(M)$ is right-angled at $M$ (resp. $O$ ), that is when $M$ is at $T$ or $U$ (resp. $V$ or $W$ ).

Lemma 2.2. When $M$ describes $\Gamma$ and $a \neq 0$, the surface of $\operatorname{Tr}(M)$ is maximum and equal to $\frac{1}{2}|a| e$ when $\gamma=\frac{\pi}{2}$.
Proof. Clear since the surface of $O M C$ is the unsigned area $\frac{1}{2}|a| e \sin \gamma$. The maximum is achieved for $M$ at $L^{+}=(a, e)$ and $L^{-}=(a,-e)$ so that $O L^{+}=O L^{-}=\sqrt{a^{2}+e^{2}}$, see Figure 2. If $a=0, \operatorname{Tr}(M)$ is degenerate.

Let $\hat{\alpha}$ denote $\angle\left(O C, O L^{+}\right)$and $\hat{\beta}=\angle\left(L^{+} C, L^{+} O\right)$, then $\tan \hat{\alpha}=\frac{e}{|a|}=\frac{1}{\tan \hat{\beta}}$. When $g^{2}>0$ (resp. < 0) the angles $\hat{\alpha}$ and $\phi$ (resp. $\hat{\beta}$ and $\psi$ ) are related by $\hat{\alpha}=\tan ^{-1} \sin \phi$ (resp. $\hat{\beta}=$ $\tan ^{-1} \sin \psi$ ).

We now suppose that $g^{2}=0$, see Figure 3 where $\Lambda=O$.


Figure 3. $\operatorname{Tr}(M)$ when $g^{2}=0$
When $|a|=e$, the triangle $O M C$ is isosceles with $\alpha=\beta$. It is right-angled if $M$ is at $L^{+}$or $L^{-}: \hat{\alpha}=\hat{\beta}=\frac{\pi}{4}(\tan \hat{\alpha}=\tan \hat{\beta}=1)$ and equilateral if $\alpha=\beta=\gamma=\frac{\pi}{3}, \delta=\frac{\pi}{6}$. It is degenerate as $O C$ when $\alpha=\beta=\frac{\pi}{2}, \gamma=0$.
2.5. Similar triangles when $\alpha$ or $\beta=\frac{\pi}{2}, g^{2} \neq 0$. We suppose that $\alpha=\phi, \beta=\frac{\pi}{2}$ for $g^{2}>0$ or $\alpha=\frac{\pi}{2}, \beta=\psi$ for $g^{2}<0$, see Figure 4. The location of the generic point $M$ is either $T$ or $U$ (resp. $V$ or $W$ ) if $g^{2}>0\left(\right.$ resp. $\left.g^{2}<0\right)$.


Figure 4. Similar right-angled triangles for $g^{2} \neq 0$
When $|a|>e$, let $P$ inside $\Gamma$ be the projection of $T$ or $U$ on the spectral axis. The triangles $O T C$ and $C P T$ are similar, with ratio $\frac{|a|}{e}>1$. Moreover $\cos \phi=\frac{g}{|a|}=\frac{O P}{g}$, hence $O P=h=\frac{g^{2}}{|a|}$, $P T=\frac{g e}{|a|}$.

When $|a|<e$, let the tangent to $\Gamma$ at $V$ intersect the real axis at $P$ outside $\Gamma$. The triangles $O V C$ and $V P C$ are also similar with ratio $\frac{e}{P C}$, with $P C=P O+|a|$ and $P O=\frac{-g^{2}}{|a|}$, hence $P C=\frac{1}{|a|}\left(-g^{2}+a^{2}\right)=\frac{e^{2}}{|a|}$ yields the value $\frac{e}{P C}=\frac{|a|}{e}<1$.

## 3. The invariant plane $\mathbf{M}$ spanned by $q$ and $q^{\prime}$

3.1. Generalities. Let $q$ and $q^{\prime}$ be an orthonormal pair of eigenvectors associated with $\lambda$ and $\lambda^{\prime}:\|q\|=\left\|q^{\prime}\right\|=1$ and $\left\langle q, q^{\prime}\right\rangle=q^{\prime H} q=0$. The subspace $\mathbf{M}$ of complex (resp. real) linear combinations of $q$ and $q^{\prime}$ is a complex (resp. real) plane invariant under the action of $A$ when
$K=\mathbb{C}($ resp. $\mathbb{R})$. Therefore $\mathbf{M}$ lives in $\mathbb{R}^{4}\left(\right.$ resp. $\left.\mathbb{R}^{2}\right)$ : the real interpretation of $\mathbf{M}$ differs whether $K=\mathbb{C}$ or $\mathbb{R}$. We suppose in this Section that $A$ is symmetric, hence $\mathbf{M}$ is isomorphic to the real plane $\mathbb{R}^{2}$. The treatment of $A$ hermitian is deferred to Sections 4 and 5 .
3.2. $K=\mathbb{R}$ : the triangle $O M^{\prime} C^{\prime}$ in the invariant 2 D -plane M . We consider the real combination $u=(\cos \theta) q+(\sin \theta) q^{\prime}$, with unit norm $\|u\|=1, \theta \in[0,2 \pi[$. When $\theta$ varies in $[0,2 \pi[$, the vector $u$ describes, in the invariant plane $\mathbf{M}$, the unit circle ( $C$ ) centered at $O$ and passing through the eigenvectors $\pm q, \pm q^{\prime}$.

We define $B=A-a I$, a matrix which commutes with $A: F=A B=B A=A^{2}-a A$.
Lemma 3.1. $\|B u\|=e$ and $\left\|F u-e^{2} u\right\|=|a| e$ for any $u \in(C)$.
Proof. $B u=\cos \theta(\lambda-a) q+\sin \theta\left(\lambda^{\prime}-a\right) q=e \tilde{u}$ with $\tilde{u}=-q \cos \theta+q^{\prime} \sin \theta,\|\tilde{u}\|=1$. Hence $\|B u\|=e$ for any $\theta$. Next $F u-e^{2} u=e a \tilde{u}$.

When $u$ is not an eigenvector $\left(\theta \notin\left\{0, \frac{\pi}{2}, \frac{3 \pi}{2}, \pi\right\}\right)$ and $a \neq 0$, the 3 vectors $a u, A u$ and $B u$ are linearly independent. Since $B u=A u-a u$ they form a non degenerate triangle $\operatorname{Tr}(u)=O M^{\prime} C^{\prime}$, see Figure 5 (a) when $a>e$ and (b) when $a<e$ which displays the circle $\Gamma^{\prime}$ centered at $C^{\prime}$ with radius $e$. In order that $\operatorname{Tr}(u)$ be non degenerate $\left(C^{\prime} \neq O\right)$, we assume below that $a \neq 0$ when $g^{2}<0$.


Figure 5. $\quad \operatorname{Tr}(u)$ as a function of $\theta$ when $K=\mathbb{R}$
The point $M^{\prime}=\left(r, r^{\prime}\right)$ in the basis $\left\{q, q^{\prime}\right\}$ is the meeting point on $\Gamma^{\prime}$ of $A u$ and $B u,\|B u\|=$ $C^{\prime} M^{\prime}=e$.
Lemma 3.2. The point $M^{\prime}=\left(r, r^{\prime}\right)$ describes the ellipse of equation $\left(\frac{r}{\lambda}\right)^{2}+\left(\frac{r^{\prime}}{\lambda^{\prime}}\right)^{2}=1$ iff $g^{2} \neq 0$. If $g^{2}=0$, the ellipse is reduced to a segment.
Proof. $A u=\lambda \cos \theta q+\lambda^{\prime} \sin \theta q^{\prime}=r q+r^{\prime} q^{\prime}$, hence $\cos ^{2} \theta+\sin ^{2} \theta=1=\left(\frac{r}{\lambda}\right)^{2}+\left(\frac{r^{\prime}}{\lambda^{\prime}}\right)^{2}$ when $\lambda \lambda^{\prime} \neq 0$. If $g^{2}=0, \lambda=0<\lambda^{\prime}, r=0$ and $r^{\prime} \in\left[-\lambda^{\prime}, \lambda^{\prime}\right]$ (say).

One can check readily on Figure 5 (a) and (b) that $\angle\left(C^{\prime} O, C^{\prime} M^{\prime}\right)=2 \angle\left(D^{\prime} O, D^{\prime} M^{\prime}\right)$ as ordinary angles in $] 0, \pi\left[\right.$, and on Figure 5 (c) that $\angle(-u, \tilde{u})=2 \theta$. The restriction $\theta \in\left[0, \frac{\pi}{2}\right]$ is sufficient to display the complete geometric evolution of $\operatorname{Tr}(u)$. Therefore we may restrict $\theta$ to describe the first quadrant for simplicity.
3.3. When $\delta=\theta$ in $] 0, \frac{\pi}{2}[$. We may compare the triangle $\operatorname{Tr}(M)=O M C$ parameterised by $\delta$ and the triangle $\operatorname{Tr}(u)=O M^{\prime} C^{\prime}$ parameterised by $\theta$.
Corollary 3.3. The equality $\delta=\theta$ in $] 0, \frac{\pi}{2}[$ yields the congruence $\operatorname{Tr}(M)=\operatorname{Tr}(u)$.

Proof. The triangles have two fixed side lengths $O C=O C^{\prime}=|a|$ and $C M=C^{\prime} M^{\prime}=e$. Each pair of sides envelops the same (ordinary) angle $\gamma$ if $\delta=\theta$ in $] 0, \frac{\pi}{2}[: \gamma=2 \theta=2 \delta \in] 0, \pi[$.

When $M^{\prime}$ describes its ellipse once in $\mathbf{M}, \theta$ varies in $[0,2 \pi[$. This entails that the corresponding point $M$ describes its circle $\Gamma$ twice in the spectral plane.
3.4. Local optimisation in M. The key Corollary 3.3 allows us to transfer the results of Section 2 about $\operatorname{Tr}(M)$ to the triangle $\operatorname{Tr}(u)$. The optimality properties $\alpha \leq \phi\left(g^{2}>0\right), \beta \leq \psi$ $\left(g^{2}<0\right)$ are valid for the corresponding ordinary angles $\alpha=\angle(u, A u), \beta=\angle(B u, A u)$ in $O M^{\prime} C^{\prime}$.

When $g^{2}>0$, we define $w_{+}$and $w_{+}^{\prime}$ as the positive square roots of $w_{+}^{2}=\frac{\lambda^{\prime}}{\lambda+\lambda^{\prime}}>0$ and $w_{+}^{\prime 2}=\frac{\lambda}{\lambda+\lambda^{\prime}}>0$. Indeed $\lambda^{\prime}\left(\lambda+\lambda^{\prime}\right)=\lambda^{\prime 2}+g^{2}$ and $\lambda\left(\lambda+\lambda^{\prime}\right)=\lambda^{2}+g^{2}$ are positive quantities. Then we define the 4 vectors $v_{+}=\varepsilon w_{+} q+\varepsilon^{\prime} w_{+}^{\prime} q^{\prime}$, with $\varepsilon= \pm 1, \varepsilon^{\prime}= \pm 1$, and $D_{+}=\left\{v_{+}\right\}$.

When $g^{2}<0$, we define analogously $w_{-}$and $w_{-}^{\prime}$ as the positive square roots of $w_{-}^{2}=\frac{\lambda^{\prime}}{\lambda^{\prime}-\lambda}$ and $w_{-}^{\prime 2}=-\frac{\lambda}{\lambda^{\prime}-\lambda}$, after we check that $\lambda^{\prime}\left(\lambda^{\prime}-\lambda\right)=\lambda^{\prime 2}-g^{2}>0$ and $-\lambda\left(\lambda^{\prime}-\lambda\right)=\lambda^{2}-g^{2}>0$. In this case we define the four vectors $v_{-}=\varepsilon w_{-} q+\varepsilon^{\prime} w_{-}^{\prime} q^{\prime}$, with $\varepsilon= \pm 1, \varepsilon^{\prime}= \pm 1$ and $D_{-}=\left\{v_{-}\right\}$.
Theorem 3.4. i) When $g^{2}>0$, the minimum value $\cos \phi=\frac{g}{|a|}$ is achieved by any $v_{+}$in $D_{+}$ and $\left\langle B v_{+}, A v_{+}\right\rangle=0,\left\|A v_{+}\right\|=\frac{1}{e}\left\|F v_{+}\right\|=g$.
ii) When $g^{2}<0$, the minimum value $\cos \psi=\frac{|g|}{e}$ is achieved by any $v_{-}$in $D_{-}$and $\left\langle v_{-}, A v_{-}\right\rangle=$ $0,\left\|A v_{-}\right\|=\frac{1}{e}\left\|F v_{-}\right\|=|g|$.
Proof. i) When $g^{2}>0$ and $v_{+}=w_{+} q+w_{+}^{\prime} q^{\prime}, A v_{+}=\lambda w_{+} q+\lambda^{\prime} w_{+}^{\prime} q^{\prime}, B v_{+}=e\left(-w_{+} q+w_{+}^{\prime} q^{\prime}\right)$ and $B A v_{+}=e\left(-\lambda w_{+} q+\lambda^{\prime} w_{+}^{\prime} q^{\prime}\right)$. Therefore $\left\|A v_{+}\right\|=g,\left\langle v_{+}, A v_{+}\right\rangle=\frac{g^{2}}{a}$, hence $\cos \angle\left(v_{+}, A v_{+}\right)=\frac{g}{a}$. Therefore $|\cos \alpha| \geq \cos \phi=\frac{g}{|a|}$. Moreover $\left\langle B v_{+}, A v_{+}\right\rangle=0$ as expected. The vector $\frac{1}{|a|} F v_{+}$is orthogonal to $v_{+}$with length $P T=P U=\frac{e g}{|a|}$.
ii) When $g^{2}<0$ and $v_{-}=w_{-} q+w_{-}^{\prime} q^{\prime}, A v_{-}=\lambda w_{-} q+\lambda^{\prime} w_{-}^{\prime} q^{\prime}$ and $\left\|A v_{-}\right\|=|g|=\sqrt{-\lambda \lambda^{\prime}}$; $B v_{-}=e\left(-w_{-} q+w_{-}^{\prime} q^{\prime}\right)$ and $\left\langle B v_{-}, A v_{-}\right\rangle=-g^{2}=|g|^{2}$. Thus $\cos \angle\left(B v_{-}, A v_{-}\right)=\frac{|g|}{e}$, that is $\cos \beta \geq \cos \psi=\frac{|g|}{e}$. Finally $\left\langle v_{-}, A v_{-}\right\rangle=0$ and $\frac{1}{|a|} F v_{-}$is orthogonal to $B v_{-}$with length $P V=P W=\frac{e|g|}{|a|}$.
3.5. Catchvectors and antieigenvectors. When $A$ is indefinite, the spectral coupling of eigenvalues with different sign: $\lambda<0<\lambda^{\prime}$ yields the existence of 4 vectors $v_{-}$with an orthogonal image $A v_{-}:\left\langle v_{-}, A v_{-}\right\rangle=0$. These vectors are the "most turned" vectors by $A$ locally in M: their image direction being orthogonal, it is the "furthest" from their own direction. Such vectors, which can exist only when $A$ is indefinite, truly deserve to be called antieigenvectors. Their dynamics under $A$ is the opposite of that for an eigenvector, whose direction is invariant under the action of $A$. To avoid ambiguity the vectors $v_{+}$in $D_{+}$when $g^{2}>0$ are called catchvectors.

Remark 3.1. Because the theory of Gustafson assumes that $A$ is positive definite, all spectral couplings are realised by positive eigenvalues with $a>0$, and $g^{2}=\lambda \lambda^{\prime}>0$ is the rule. Therefore $\cos \alpha$ is bounded from below by $\frac{g}{a}>0$. Because we put no other assumption on $A \in \mathbb{R}^{n \times n}$ than symmetry, we are led to depart from Gustafson's nomenclature. When $A$ is indefinite, it becomes necessary to distinguish between catch- and antieigen- vectors.

But it turns out that our work differs from Gustafson's theory on some more fundamental grounds. Because of the basic assumption that $A$ is symmetric positive definite, his theory focuses mainly on the specific angle $\alpha=\angle(u, A u)$ obtained for ( $\lambda_{1}, \lambda_{n}$ ); it ignores the key role of the triangles $\operatorname{Tr}(M)$ in the spectral plane $\boldsymbol{\Sigma}$ and $\operatorname{Tr}(u)$ in the invariant plane $\mathbf{M} \subset \mathbb{R}^{n}$, and of the information brought by their three angles $\alpha, \beta, \gamma=\frac{\theta}{2}=\frac{\delta}{2}$. The full extent to which we have developed Gustafson's seminal ideas will be unravelled as we proceed.
3.6. Middle vectors. The previous section has dealt with the orthogonality $\beta=\frac{\pi}{2}$ when $g^{2}>0$ and $\alpha=\frac{\pi}{2}$ when $g^{2}<0$. We turn to the third angle $\gamma=2 \theta$. We define $\hat{w}=\frac{1}{\sqrt{2}}$ and $\hat{v}=\hat{w}\left(\varepsilon q+\varepsilon^{\prime} q^{\prime}\right), \hat{D}=\left\{\hat{v}, \varepsilon\right.$ and $\left.\varepsilon^{\prime}= \pm 1\right\}$.

Lemma 3.5. The minimum value $\langle a u, B u\rangle=0$ is achieved, for $a \neq 0$, at $\hat{v}=\frac{1}{\sqrt{2}}\left(q+q^{\prime}\right)$.
Proof. $\langle a u, B u\rangle=a e\langle u, \tilde{u}\rangle=0 \Leftrightarrow \theta=\frac{\pi}{4} \Leftrightarrow \cos \theta=\sin \theta=\frac{1}{\sqrt{2}} \Leftrightarrow \gamma=\frac{\pi}{2}$.
Theorem 3.6. When $a \neq 0$, the 4 triangles $\operatorname{Tr}(\hat{v}), \hat{v} \in \hat{D}$, have the maximal surface $\frac{1}{2}|a| e$.
Proof. The surface of the triangle $\operatorname{Tr}(u)$ is $\Sigma(u)=\frac{1}{2}|a| e \sin \gamma$ which achieves its maximum for $u=\hat{v} \Leftrightarrow \gamma=\frac{\pi}{2}$.

The quantity $2 \Sigma(u)$ is called the influence of $u$. It is nonzero when $\operatorname{Tr}(u)$ is non degenerate ( $u \notin\left\{ \pm q, \pm q^{\prime}\right\}$ or $a \neq 0$ ).

When comparing $D_{ \pm}$and $\hat{D}$, we observe that the vectors $\hat{v}$ are independent of the values $\lambda<\lambda^{\prime}$. These vectors are called middle vectors since they are the bisectors of the eigenvectors. They have the largest influence $|a| e$, which is the surface of any of the two rectangles with diagonal $O L^{+}$or $O L^{-}$, see Figure 2. Moreover Lemma 3.1 tells us that the maximal surface $|a| e$ is precisely the norm of $F u-e^{2} u$ for any $u \in(C)$.

Remark 3.2. The generic concept of a middle vector is absent from Gustafson's theory. It only appears in a statistical setting under the guise of an "inefficient" vector, see Section 6.2 and [Gustafson, 2012, p. 190].
3.7. $A$ is positive definite. Let $A$ be symmetric positive definite with eigenvalues $0<\lambda_{1} \leq$ $\ldots \leq \lambda_{n}$. Choosing the extreme pair $\left(\lambda_{1}, \lambda_{n}\right)$ we set $a_{*}=\frac{\lambda_{1}+\lambda_{n}}{2}>0, e_{*}=\frac{\lambda_{n}-\lambda_{1}}{2}, g_{*}=\sqrt{\lambda_{1} \lambda_{n}}$.

Lemma 3.7. For any coupling $\lambda<\lambda^{\prime}$ which is not extreme the following inequalities hold: $\frac{e}{a}<\frac{e_{*}}{a_{*}}<1 \Leftrightarrow \frac{g}{a}>\frac{g_{*}}{a_{*}}>0, a e<a_{*} e_{*}$.
Proof. First, $\frac{e}{a}<\frac{e_{*}}{a_{*}}<1 \Leftrightarrow \lambda_{1} \lambda\left(\frac{\lambda_{n}}{\lambda_{1}}-\frac{\lambda^{\prime}}{\lambda}\right)>0$. Since $\frac{g}{a}=\frac{\sqrt{a^{2}-e^{2}}}{a}=\sqrt{1-\left(\frac{e}{a}\right)^{2}}$, we get $\frac{g}{a}>\frac{g_{*}}{a_{*}}>0$. Second, $a<a_{*}$ and $e<e_{*}$ entail $a e<a_{*} e_{*}$.

Corollary 3.8. The maximum turning angle $\phi(A)$ satisfies $\cos \phi(A)=\frac{g_{*}}{a_{*}}, \sin \phi(A)=\frac{e_{*}}{a_{*}}$. The associated middle vectors have the highest influence $a_{*} e_{*}=\frac{1}{4}\left(\lambda_{n}^{2}-\lambda_{1}^{2}\right)$.

Proof. Direct consequence of Lemma 3.7. The weights $w_{+, *}, w_{+, *}^{\prime}$ of the catchvectors $v_{+, *}$ satisfy $\frac{w_{+, *}}{w_{+, *}^{\prime}}=\sqrt{\frac{\lambda_{n}}{\lambda_{1}}}=\operatorname{cond}\left(A^{1 / 2}\right)$, where $\operatorname{cond}(A)=\sup _{\lambda^{\prime}>\lambda} \frac{\lambda^{\prime}}{\lambda}$ is the condition number of $A$ expressed with the euclidean norm $\|\cdot\|$. The middle vectors are the bisectors of any two orthonormal eigenvectors $q_{*}, q_{*}^{\prime}$ for $\lambda_{1}$ and $\lambda_{n}: \hat{v}_{*}=\frac{1}{\sqrt{2}}\left(\varepsilon q_{*}+\varepsilon^{\prime} q_{*}^{\prime}\right), \varepsilon$ and $\varepsilon^{\prime}= \pm 1$.

For future reference we denote $\gamma(A)$ the trigonometric complement of $\phi(A): \phi(A)+\gamma(A)=$ $\frac{\pi}{2}=\beta(A)$.
3.8. A difference between $\operatorname{Tr}(M)$ and $\operatorname{Tr}(u)$. Despite their congruence when $a \neq 0$, the triangles $\operatorname{Tr}(M)$ and $\operatorname{Tr}(u)$ do not process spectral information in an identical way. In the spectral plane $\boldsymbol{\Sigma}$, the axis $O C$ is fixed whereas in $\mathbf{M}$, only $O$ is fixed: $C^{\prime}$ describes the circle $|a|(C)$.

We saw in Section 2 that $\operatorname{Tr}(M)$ can be right-angled at the vertex $M\left(g^{2}>0\right)$ or $O\left(g^{2}<0\right)$ and $C\left(g^{2} \neq 0\right)$ for $M$ at two locations on $\Gamma$, namely $(T, U)$ or $(V, W)$ and $\left(L^{+}, L^{-}\right)$. By comparison $\operatorname{Tr}(u)$ can be right-angled at $M^{\prime}\left(g^{2}>0\right)$ or $O\left(g^{2}<0\right)$ and $C^{\prime}\left(g^{2} \neq 0\right)$ when $u$ belongs to the sets $D_{+}$or $D_{-}$and $\hat{D}$. And each set consists of four vectors leading to eight right-angled triangles ( 2 triangles per vector).

## 4. Variational principles

In this Section, $A$ is an arbitrary hermitian matrix, $A=A^{H} \in \mathbb{C}^{n \times n}$ where the ground field is $K=\mathbb{C}$ which can be specialised to be $\mathbb{R}$ ( $A$ symmetric).

For any hermitian matrix $Y$, the ratio $\frac{x^{H} Y x}{\|x\|\|Y x\|}$ is real in $[-1,1]$ for any $0 \neq x \in \mathbb{C}^{n}$. When $x \in \mathbb{R}^{n}$, the ratio can be interpreted as $\cos \measuredangle(x, Y x)$ thanks to Cauchy's inequality. Because $x$ defines a real direction, the angle $\mathcal{Y}(x)=\measuredangle(x, Y x)$ is a direction (or rotation) angle defined $\bmod 2 \pi$ between the directions spanned by the real unit vectors $\frac{x}{\|x\|}$ and $y=\frac{Y x}{\|Y x\|}$. Such a geometric interpretation is not readily available for $A$ hermitian since $x \in \mathbb{C}^{n}$. In particular the number Arcos $\frac{|\langle x, y\rangle|}{\|x\|\|y\|}$ in $\left[0, \frac{\pi}{2}\right]$ which is commonly referred to as angle $(x, y)$ is of an analytic, rather than geometric, nature. The question is discussed further in Section 5.1.

To avoid any ambiguity we use two distinct notations to represent an "angle" according to $K$ :

- $K=\mathbb{R}^{n}, \mathcal{Y}(x)=\measuredangle(x, Y x) \in[0,2 \pi]$ with geometric and analytic meaning,
- $K=\mathbb{C}^{n}$, angle $(x, y) \in\left[0, \frac{\pi}{2}\right]$ with analytic meaning only.
4.1. A preparatory lemma (M.M. Rincon-Camacho). Let be given two hermitian matrices $Y$ and $Z$, the product $Y Z$ is hermitian iff $Y$ and $Z$ commute. We consider the real functional

$$
\begin{equation*}
c(x)=\frac{x^{H} Y Z x}{\|Y x\|\|Z x\|} \in \mathbb{R}, \quad 0 \neq x \in K^{n} \backslash(\operatorname{Ker} Y \cup \operatorname{Ker} Z) \tag{4.1}
\end{equation*}
$$

where $Y$ and $Z$ are hermitian and commute: $Y Z=Z Y$. Thus $|c(x)|=\cos (\operatorname{angle}(Y x, Z x))$ if $K=\mathbb{C}$ or $c(x)=\cos \measuredangle(Y x, Z x)$ if $K=\mathbb{R}$.

Lemma 4.1. The Euler equation for (4.1) is given for $0 \neq x \in K^{n} \backslash(\operatorname{Ker} Y \cup \operatorname{Ker} Z)$, $\langle Y x, Z x\rangle \neq 0$ by:

$$
\begin{equation*}
\frac{Y^{2} x}{\|Y x\|^{2}}-\frac{2 Y Z x}{\langle Y x, Z x\rangle}+\frac{Z^{2} x}{\|Z x\|^{2}}=0 \tag{4.2}
\end{equation*}
$$

Proof. When $K=\mathbb{R}$ and $Y=I, Z=A$ symmetric positive definite, the proof is easily adapted from that of Theorem 3.2 on p. 36 in [Gustafson, 2012]. For the sake of completeness we adapt below the proof to the general case $Y Z$ hermitian, $K=\mathbb{C}$.

In order to find (4.2), one looks for those $x$ in $\mathbb{C}^{n} \backslash\{0\}$ which make the directional derivative

$$
\frac{d c(x)}{d y}(\varepsilon=0)=\lim _{\varepsilon \rightarrow 0} \frac{1}{\varepsilon}(c(x+\varepsilon y)-c(x)), \quad \varepsilon \in \mathbb{C}, \quad 0 \neq y \in \mathbb{C}^{n}
$$

vanish for all directions $y \in \mathbb{C}^{n} \backslash\{0\}$. We consider for $\varepsilon>0$ small enough

$$
c(x+\varepsilon y)-c(x)=\frac{(x+\varepsilon y)^{H} Y Z(x+\varepsilon y)}{\|Y(x+\varepsilon y)\|\|Z(x+\varepsilon y)\|}-\frac{x^{H} Y Z x}{\|Y x\|\|Z x\|}=\frac{N}{D}, \quad x \notin \operatorname{Ker} Y \cup \operatorname{Ker} Z
$$

with

$$
\begin{aligned}
N= & \left(\langle Y Z x, x\rangle+2 \mathfrak{R} \varepsilon\langle Y Z y, x\rangle+|\varepsilon|^{2}\langle Y Z y, y\rangle\right)\|Y x\|\|Z x\| \\
& -\langle Y Z x, x\rangle\left(\|Y x\|^{2}+2 \mathfrak{R} \varepsilon\langle Y x, Y y\rangle+|\varepsilon|^{2}\|Y y\|^{2}\right)^{1 / 2}\left(\|Z x\|^{2}+2 \mathfrak{R} \varepsilon\langle Z x, Z y\rangle+|\varepsilon|^{2}\|Z y\|^{2}\right)^{1 / 2} \\
D= & \left(\|Y x\|^{2}+2 \mathfrak{R} \varepsilon\langle Y x, Y y\rangle+|\varepsilon|^{2}\|Y y\|^{2}\right)^{1 / 2}\left(\|Z x\|^{2}+2 \mathfrak{R}\langle Z x, Z y\rangle+|\varepsilon|^{2}\|Z y\|^{2}\right)^{1 / 2}\|Y x\|\|Z x\| .
\end{aligned}
$$

Clearly $D \rightarrow(\|Y x\|\|Z x\|)^{2}$ as $\varepsilon \rightarrow 0$. In order to find $\lim _{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} N$, we consider limited series expansions in $\varepsilon$ for the terms in $N$ to be subtracted:

1) $\|Y(x+\varepsilon y)\|=\left(\|Y x\|^{2}+f(\varepsilon)\right)^{1 / 2}=\|Y x\|+\frac{1}{2} \frac{1}{\|Y x\|} f(\varepsilon)-\frac{1}{8} \frac{1}{\|Y x\|^{3}} f^{2}(\varepsilon)+\ldots=\|Y x\|+r(\varepsilon)$,
2) $\|Z(x+\varepsilon y)\|=\left(\|Z x\|^{2}+g(\varepsilon)\right)^{1 / 2}=\|Z x\|+\frac{1}{2} \frac{1}{\|Z x\|} g(\varepsilon)-\frac{1}{8} \frac{1}{\|Z x\|^{3}} g^{2}(\varepsilon)+\ldots=\|Z x\|+t(\varepsilon)$.

Here $f(\varepsilon)=2 \mathfrak{R} \varepsilon\langle Y x, Y y\rangle+|\varepsilon|^{2}\|Y y\|^{2}$ and $g(\varepsilon)=2 \mathfrak{R} \varepsilon\langle Z x, Z y\rangle+|\varepsilon|^{2}\|Z y\|^{2}$ are functions which depend on $\varepsilon$ taken sufficiently small relative to $\|Y x\|$ and $\|Z x\|$ respectively. Thus

$$
\begin{aligned}
N= & \left(\langle Y Z x, x\rangle+2 \mathfrak{R} \varepsilon\langle Y Z s, x\rangle+|\varepsilon|^{2}\langle Y Z s, s\rangle\right)\|A x\|\|Z x\| \\
& -\langle Y Z x, x\rangle(\|Y x\|+r(\varepsilon))(\|Z x\|+t(\varepsilon)) \\
= & \left(2 \mathfrak{R} \varepsilon\langle Y Z y, x\rangle+|\varepsilon|^{2}\langle Y Z y, y\rangle\right)\|Y x\|\|Z x\|-\langle Y Z x, x\rangle(\|Y t\| t(\varepsilon)+\|Z x\| r(\varepsilon)) \\
= & \left(2 \mathfrak{R} \varepsilon\langle Y Z y, x\rangle+|\varepsilon|^{2}\langle Y Z y, y\rangle\right)\|Y x\|\|Z x\|-\langle Y Z x, x\rangle\left(\frac{1}{2} \frac{\|Y x\|}{\|Z x\|} g(\varepsilon)+\frac{1}{2}\|Z x\| x \|\right. \\
= & \left(2 \mathfrak{R} \varepsilon\langle Y Z y, x\rangle+|\varepsilon|^{2}\langle Y Z y, y\rangle\right)\|Y x\|\|Z x\| \\
& -\langle Y Z x, x\rangle\left[\frac{\|Y x\|}{\|Z x\|}\left(\mathfrak{R} \varepsilon\langle Z x, Z y\rangle+\frac{1}{2}|\varepsilon|^{2}\|Z y\|^{2}\right)+\frac{\|Z x\|}{\|Y x\|}\left(\mathfrak{R} \varepsilon\langle Y x, Y y\rangle+\frac{1}{2}|\varepsilon|^{2}\|Y y\|^{2}\right)\right]
\end{aligned}
$$

and

$$
\begin{aligned}
\frac{N}{\varepsilon}= & (2 \mathfrak{R}\langle Y Z y, x\rangle+\bar{\varepsilon}\langle Y Z y, y\rangle)\|Y x\|\|Z x\| \\
& -\langle Y Z x, x\rangle\left[\frac{\|Y x\|}{\|Z x\|}\left(\mathfrak{R}\langle Z x, Z y\rangle+\frac{1}{2} \bar{\varepsilon}\|Z y\|^{2}\right)+\frac{\|Z x\|}{\|Y x\|}\left(\mathfrak{R}\langle Y x, Y y\rangle+\frac{1}{2} \bar{\varepsilon}\|Y y\|^{2}\right)\right] .
\end{aligned}
$$

Finally we have

$$
\begin{aligned}
\lim _{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} \frac{N}{D} & =\frac{2 \mathfrak{R}\langle Y Z y, x\rangle}{\|Y x\|\|Z x\|}-\langle Y Z x, x\rangle\left(\frac{\mathfrak{R}\langle Z x, Z y\rangle}{\|Y x\|\|Z x\|^{3}}+\frac{\mathfrak{R}\langle Y x, Y y\rangle}{\|Y x\|^{3}\|Z x\|}\right) \\
& =\frac{1}{\|Y x\|\|Z x\|} \mathfrak{R}\langle y, V\rangle, \quad \text { where } V=2 Y Z x-\langle Y Z x, x\rangle\left(\frac{Z^{2} x}{\|Z x\|^{2}}+\frac{Y^{2} x}{\|Y x\|^{2}}\right) \in \mathbb{C}^{n} .
\end{aligned}
$$

The variational calculus imposes that $\mathfrak{R}\langle y, V\rangle=0$ for any $0 \neq y \in \mathbb{C}^{n}$. Indeed, if $V^{H} y=i b$, then $V^{H}(i y)=-b \in \mathbb{R}$ should also be 0 . Therefore $\mathfrak{R}\langle y, V\rangle=0$ for all $y \neq 0 \Leftrightarrow V=0 \Leftrightarrow x$ satisfies (4.2).
4.2. The catchvectors: $Y=I, Z=A$. When we choose $Z=A$ and $Y=I$, the Euler equation (4.2) becomes

$$
\begin{equation*}
A^{2} x-2 \frac{\|A x\|^{2}}{\langle x, A x\rangle} A x+\left(\frac{\|A x\|}{\|x\|}\right)^{2} x=0 \tag{4.3}
\end{equation*}
$$

for $\langle x, A x\rangle \neq 0, x \in \mathbb{C}^{n} \backslash \operatorname{Ker} A$.
In order to solve (4.3) we set $\frac{\|A x\|^{2}}{\langle x, A x\rangle}=k(x)=k$ and $\left(\frac{\|A x\|}{\|x\|}\right)^{2}=l(x)=l>0$. Then with $A=Q D Q^{H}, y=Q^{H} x,\|y\|=\|x\|$, (4.3) can be written

$$
\left(D^{2}-2 k D+l I\right) y=0, \quad y=\left(y_{i}\right) \in \mathbb{C}^{n},
$$

that is, with $D=\operatorname{diag}\left(\mu_{i}\right)$ :

$$
\begin{equation*}
\left(\mu_{i}^{2}-2 k \mu_{i}+l\right) y_{i}=0, \quad i=1, \ldots, n . \tag{4.4}
\end{equation*}
$$

Observe that the discriminant is $k^{2}-l=\|A x\|^{2}\left(\frac{\|A x\|^{2}}{\langle x, A x\rangle^{2}}-\frac{1}{\|x\|^{2}}\right)$ : its sign is that of $(\|x\|\|A x\|-$ $|\langle x, A x\rangle|) \geq 0$ by Cauchy's inequality. Eq. (4.3) is obviously satisfied when $x$ is an eigenvector: $A x=\lambda x$. Let us assume that $x$ and $A x$ are independent. Eq. (4.4) entails $y_{j}=0$ for all $j$ such that $\mu_{i}^{2}-2 k \mu_{i}+l \neq 0$ and vice versa. We denote $\lambda<\lambda^{\prime}$ the two distinct roots $\left\{\mu_{j}, \mu_{j}^{\prime}\right\}$ of the quadratic equation. The corresponding multiplicities are $m, m^{\prime}$ and the associated eigensubspaces are $\mathcal{M}, \mathcal{M}^{\prime}$ with $\operatorname{dim} \mathcal{M}=m, \operatorname{dim} \mathcal{M}^{\prime}=m^{\prime}$. As before, $a=\frac{\lambda+\lambda^{\prime}}{2}, g^{2}=\lambda \lambda^{\prime}$.
Lemma 4.2. Let $X, X^{\prime}$ be arbitrary normalised eigenvectors in $\mathcal{M}, \mathcal{M}^{\prime}$. Any vector $x=X+X^{\prime}$ in $\mathcal{M} \oplus \mathcal{M}^{\prime},\|x\|=\sqrt{2}$, is a solution of (4.3) iff $k(x)=a$ and $l(x)=g^{2}, 0 \leq g^{2}<a^{2}$.

Proof. By definition, $\lambda=\mu_{\jmath}, \jmath \in J_{m}=\{1, \ldots, m\}$ and $\lambda^{\prime}=\mu_{\jmath^{\prime}}, \jmath^{\prime} \in J_{m^{\prime}}=\left\{1, \ldots, m^{\prime}\right\}$, and the pair $\left\{\mu_{\jmath}, \mu_{\jmath^{\prime}}\right\}$ solves the quadratic equation $r^{2}-2 a r+g^{2}=\left(r-\mu_{\jmath}\right)\left(r-\mu_{\jmath^{\prime}}\right)=0$. It follows that for $k=a$ and $l=g^{2}>0$, Eq. (4.4) is satisfied for $i \in J_{m} \cup J_{m^{\prime}}$, letting $y_{i} \in \mathbb{C}$ be arbitrary. For $i$ outside $J_{m} \cup J_{m^{\prime}}$, it is clear that $\mu_{i}^{2}-2 k \mu_{i}+l \neq 0$, imposing that $y_{i}=0$. Let $q_{i}$ be the columns of $Q=\left[q_{1}, \ldots, q_{n}\right]$ which are eigenvectors for $A$. Then $x=Q y=\sum_{\jmath \in J_{m}} y_{\jmath} q_{\jmath}+\sum_{\jmath^{\prime} \in J_{m^{\prime}}} y_{\jmath^{\prime}} q_{\jmath^{\prime}}=X+X^{\prime}$.

Let $q, q^{\prime}$ be two orthonormal eigenvectors associated with $\lambda<\lambda^{\prime}$ which span the invariant subspace $\mathbf{M}$ with 4 real dimensions when $K=\mathbb{C}$. The unit sphere $(S)$ in $\mathbb{R}^{4}$ passing through $q$ and $q^{\prime}$ consists of vectors $u=z q+z^{\prime} q^{\prime},|z|^{2}+\left|z^{\prime}\right|^{2}=1$. When are the conditions $k(u)=a$, $l(u)=g^{2}>0$ satisfied for $u \in(S)$ ?
Proposition 4.3. The solutions of Euler's equation (4.3) which are not eigenvectors are the catchvectors $v_{+}=e^{i \xi} w_{+} q+e^{i \xi^{\prime}} w_{+}^{\prime} q^{\prime}, \xi, \xi^{\prime} \in\left[0,2 \pi\left[\right.\right.$ corresponding to all couplings $\lambda<\lambda^{\prime}$ such that $g^{2}=\lambda \lambda^{\prime}>0$. They yield the critical value $c\left(v_{+}\right)=\frac{g}{a}=\operatorname{sgn}(a) \cos \phi$.
Proof. For $u=z q+z^{\prime} q^{\prime},\|u\|=1$, set $|z|^{2}=\tau,\left|z^{\prime}\right|^{2}=1-\tau$. $k(u)=a \Leftrightarrow 2 \frac{\lambda^{2} \tau+\lambda^{\prime 2}(1-\tau)}{\lambda \tau+\lambda^{\prime}(1-\tau)}=\lambda+\lambda^{\prime}$ entails $\tau=\frac{\lambda^{\prime}}{\lambda+\lambda^{\prime}}=w_{+}^{2}$ and $1-\tau=\frac{\lambda}{\lambda+\lambda^{\prime}}=w_{+}^{\prime 2}$, hence $u=v_{+}$. One checks that $l\left(v_{+}\right)=g^{2}>0$. The conclusion follows from $c\left(v_{+}\right)=\frac{\left\langle v_{+}, A v_{+}\right\rangle}{\left\|A v_{+}\right\|},\left\langle v_{+}, A v_{+}\right\rangle=\frac{g^{2}}{a}=h$. Thus $\left|c\left(v_{+}\right)\right|=\cos \phi$.

Corollary 4.4. If $A$ is positive definite, the extreme coupling $\left\{\lambda_{1}=\lambda_{\min }, \lambda_{n}=\lambda_{\max }\right\}$ yields the global minimum (resp. maximum) $\cos \phi(A)=\frac{2 \sqrt{\lambda_{1} \lambda_{n}}}{\lambda_{1}+\lambda_{n}}\left(\right.$ resp. $\sin \phi(A)=\frac{\lambda_{n}-\lambda_{1}}{\lambda_{n}+\lambda_{1}}$ ).
Proof. 1) The case $K=\mathbb{R}$ is clear. Each of the 4 catchvectors $v_{+, *}=\varepsilon w_{+, *} q+\varepsilon^{\prime} w_{+, *}^{\prime} q^{\prime}$ is rotated by $A$ by the maximal turning angle $\phi(A)$ which shows in $\operatorname{Tr}\left(v_{+, *}\right)$ and equals $\mathcal{A}(\hat{v})$. This is one of the major theorems in Gustafson's approach. Diverse applications are presented in [Gustafson, 2012, chapters 4 to 8].
2) The proof is also clear when $K=\mathbb{C}$ but the interpretation differs. The global minimum for $c(x)$ equals $\frac{g_{*}}{a_{*}}=c\left(v_{+, *}\right)$, a value which, coincidentally, equals $\cos \phi(A)$ displayed in $\boldsymbol{\Sigma}$ by $\operatorname{Tr}(T)$ and $\operatorname{Tr}(U)$ with the choice $\left(\lambda_{1}, \lambda_{n}\right)$.
4.3. The antieigenvectors: $Z=A$ indefinite, $Y=A-a I=B$. When $A$ is indefinite, Section 3 with $K=\mathbb{R}$ suggests to choose $Z=A, Y=B$, so that $c(x)=\cos \measuredangle(B x, A x)$ reduces to $\cos \beta$ when $x$ belongs to $\mathbf{M}$. The Euler equation becomes

$$
\begin{equation*}
A^{2} x-2 \frac{\|A x\|^{2}}{\langle B x, A x\rangle} A B x+\left(\frac{\|A x\|}{\|B x\|}\right)^{2} B^{2} x=0 \tag{4.5}
\end{equation*}
$$

for $\langle B x, A x\rangle \neq 0, x \in \mathbb{C}^{n} \backslash($ Ker $A \cup \operatorname{Ker} B)$. We set $k_{1}(x)=\frac{\|A x\|^{2}}{\langle B x, A x\rangle}, l(x)=\left(\frac{\|A x\|}{\|x\|}\right)^{2}$ unchanged. Using $B=A-a I$, we rewrite (4.5) under the equivalent form $A^{2} x-2 K A x+L x=0$ where $K=\frac{l-k_{1}}{1-2 k_{1}+l}$ and $L=\frac{l a^{2}}{1-2 k_{1}+l}$ depend on $x$ by means of $k_{1}(x)$ and $l(x)$ which are assumed to satisfy $2 k_{1} \neq 1+l$.
Proposition 4.5. When $A$ is indefinite, the solutions of Euler's equation (4.5) which are not eigenvectors are the antieigenvectors $v_{-}=e^{i \xi} w_{-} q+e^{i \xi^{\prime}} w_{-}^{\prime} q^{\prime}, \xi, \xi \in[0,2 \pi[$ corresponding to all couplings $\left\{\lambda, \lambda^{\prime}\right\}$ such that $\lambda<0<\lambda^{\prime}, g^{2}<0$. They yield the critical value $c\left(v_{-}\right)=\cos \psi$.
Proof. 1) Let $u=e^{i \xi} q, A u=\lambda u, B u=-e u,\langle B u, A u\rangle=-\lambda e$ and (4.5) is obviously satisfied: $\left(\lambda^{2}-2 \lambda^{2}+\lambda^{2}\right) q=0$.
2) When $u \in(S)$ is not an eigenvector, $u$ may satisfy (4.5) iff $K(u)=a$ and $L(u)=g^{2}$. When $g^{2}<0$, it is easy to check that $u=z q+z^{\prime} q^{\prime}$ should be such that $\tau=|z|^{2}=\frac{\lambda^{\prime}}{\lambda^{\prime}-\lambda}=w_{-}^{2}$. Therefore $u$ is any of the antieigenvectos $v_{-}$. Direct computation shows that the antieigenvectors are the only solutions which are not eigenvectors.
3) $c\left(v_{-}\right)=\frac{v_{-}^{H} B A v_{-}}{\left\|B v_{-}\right\|\left\|A v_{-}\right\|}=-\frac{g^{2}}{e|g|}=\frac{|g|}{e}=\cos \psi>0$ for $a \neq 0$.

When, $K=\mathbb{C}$, we still denote $D_{+}$(resp. $D_{-}$) the set of catchvectors (resp. antieigenvectors) whose cardinality is now that of the continuum. It is clear that Theorem 3.4 remains valid in a weak form. If only expresses a coincidence between $\left|c\left(v_{+}\right)\right|$and $\cos \phi$ or $c\left(v_{-}\right)$and $\cos \psi$.
4.4. The middle vectors. The choice $Y=-I, a \neq 0, Z=B=A-a I$ and the corresponding Euler equation cannot yield the middle vectors $\hat{v}=\frac{1}{\sqrt{2}}\left(e^{i \xi} q+e^{i \xi^{\prime}} q^{\prime}\right), \xi, \xi^{\prime} \in[0,2 \pi[$ because (4.2) is not defined for $x=\hat{v}$ since $\langle\hat{v}, B \hat{v}\rangle=\frac{e}{2}\left\langle e^{i \xi} q+e^{i \xi^{\prime}} q^{\prime},-e^{i \xi} q+e^{i \xi^{\prime}} q^{\prime}\right\rangle=0$. It turns out that these vectors can be characterised in another way when $A$ is invertible. Indeed, the following result holds when $\operatorname{det} A \neq 0$ :

Proposition 4.6. When $A$ is invertible, the solutions $x \in K^{n}$, distinct from eigenvectors, of the equation

$$
\begin{equation*}
A^{2} x-2 \frac{\langle x, A x\rangle}{\|x\|^{2}} A x+\frac{\langle x, A x\rangle}{\left\langle x, A^{-1} x\right\rangle} x=0 \tag{4.6}
\end{equation*}
$$

are middle vectors $\hat{v}$ associated with all couplings $\lambda<\lambda^{\prime}$ such that $\lambda \neq-\lambda^{\prime} \Leftrightarrow a \neq 0$.
Proof. 1) It is easily checked that $\langle\hat{v}, A \hat{v}\rangle=\frac{1}{2}\left(\lambda+\lambda^{\prime}\right)=a$. Because $\lambda \lambda^{\prime} \neq 0, A^{-1} \hat{v}=\frac{1}{\sqrt{2}}\left(\frac{1}{\lambda} e^{i \xi} q+\right.$ $\left.\frac{1}{\lambda^{\prime}} e^{i \xi^{\prime}} q^{\prime}\right)$ is well-defined, thus $\frac{\langle\hat{v}, A \hat{v}\rangle}{\left\langle\hat{v}, A^{-1} \hat{v}\right\rangle}=\frac{a}{a} g^{2}=g^{2}$ when $a \neq 0$.
2) Direct computation shows that the middle vectors $\hat{v}$ are the only solutions in $(S)$ which are not eigenvectors.

The equation (4.6) is called the balance equation since its (non eigenvector) solutions are the middle vectors $\hat{v}$ independently of the nonzero values $\lambda$ and $\lambda^{\prime} \neq-\lambda$.

It is noteworthy that the balance equation is defined only if $A^{-1}$ exists. However $\hat{v}$ is welldefined for any $g^{2}$ in $]-e^{2}, a^{2}[$, which includes 0 for $a \neq 0$, hence $\hat{v}$ satisfies $\langle a \hat{v}, B \hat{v}\rangle=0$. Moreover the balance equation, which provides the middle vectors, does not indicate how one could extend from $K=\mathbb{R}$ to $\mathbb{C}$, the property that $\Sigma(\hat{v})=\max _{u \in(C)} \Sigma(u)=\frac{1}{2}|a| e$ since we have no more triangles at hand.

Again, a weak form of Theorem 3.6 remains valid when $K=\mathbb{C}$.
Proposition 4.7. Given $\lambda<\lambda^{\prime}$, any middle vector in $\hat{D}$ contains the spectral information which defines in $\boldsymbol{\Sigma}$ the triangles $\operatorname{Tr}\left(L^{+}\right)$and $\operatorname{Tr}\left(L^{-}\right)$with maximal surface $\frac{1}{2}|a| e$.

And this is but one of the many properties of the middle vectors.
4.5. $A$ is positive definite: $Y=A^{1 / 2}, Z=A^{-1 / 2}$. Interestingly, the middle vectors can be indirectly related to the preparatory Lemma 4.1 when $A$ is definite thanks to the square root $Y=A^{1 / 2}$ (resp. $(-A)^{1 / 2}$ ) when $A$ is positive (resp. negative) definite. We suppose below that $A$ is positive definite, so that $\langle x, A x\rangle=\left\|A^{1 / 2} x\right\|^{2}>0$ for $x \neq 0$. Then Eq. (4.2) takes the form

$$
\begin{equation*}
\frac{A x}{\langle x, A x\rangle}-2 \frac{x}{\|x\|^{2}}+\frac{A^{-1} x}{\left\langle x, A^{-1} x\right\rangle}=0 . \tag{4.7}
\end{equation*}
$$

Multiplying (4.7) by $A$, we rewrite it in the equivalent form

$$
A^{2} x-2 \frac{\langle x, A x\rangle}{\|x\|^{2}} A x+\frac{\langle x, A x\rangle}{\left\langle x, A^{-1} x\right\rangle} x=0
$$

which is precisely the balance equation (4.6) of Proposition 4.6. The corresponding functional is $c(x)=\frac{\|x\|^{2}}{\|x\|_{A}\|x\|_{A^{-1}}}$ where $\|x\|_{A}=\langle x, A x\rangle^{1 / 2}$ denotes the elliptic norm defined by $A$.

We just have proved the
Proposition 4.8. When $A$ is positive definite, (4.6) is the Euler equation associated with $\min _{x \neq 0} \frac{\|x\|^{2}}{\|x\|_{A}\|x\|_{A^{-1}}}$ which represents either

$$
\min _{0 \neq x \in \mathbb{R}^{n}} \cos \measuredangle\left(A^{-1 / 2} x, A^{1 / 2} x\right) \text { or } \min _{0 \neq x \in \mathbb{C}^{n}} \cos \operatorname{angle}\left(A^{-1 / 2} x, A^{1 / 2} x\right)
$$

The global minimum is achieved for the extreme pair $\left\{\lambda_{1}=\lambda_{\min }, \lambda_{n}=\lambda_{\max }\right\}$ and a pair of associated middle vectors at the value $\frac{2 \sqrt{\lambda_{1} \lambda_{n}}}{\lambda_{1}+\lambda_{n}}=\cos \phi(A)$.

Moreover, there exists a remarkable relationship between middle vectors $\hat{v}$ and catchvectors $v_{+}$due to the second author.

Corollary 4.9. Let $\hat{v}$ be a middle vector, and $h=\frac{g^{2}}{a}$ be the harmonic mean. Then $v_{+}=$ $\sqrt{h} A^{-1 / 2} \hat{v}$ is a catchvector, $A v_{+}=\sqrt{h} A^{1 / 2} \hat{v}$ and $\left\langle A^{-1 / 2} \hat{v}, A^{1 / 2} \hat{v}\right\rangle=\frac{1}{h}\left\langle v_{+}, A v_{+}\right\rangle=1=\|\hat{v}\|^{2}$.

Proof. [Rincon-Camacho, 2015b]. By straightforward calculation: $\left\langle v_{+}, A v_{+}\right\rangle=h$.
If $A$ is symmetric $(K=\mathbb{R})$, direction angles are well-defined. Since $\cos \measuredangle\left(A^{1 / 2} \hat{v}, \hat{v}\right)=$ $\cos \measuredangle\left(A^{-1 / 2} \hat{v}, \hat{v}\right)=\frac{\sqrt{\lambda^{\prime}}+\sqrt{\lambda}}{2 \sqrt{a}}$, then $\measuredangle\left(v_{+}, \hat{v}\right)=\measuredangle\left(\hat{v}, A v_{+}\right)$and $v_{+}, F_{0} v_{+}$are orthogonal $\left(\beta=\frac{\pi}{2}\right)$, the middle vector $\hat{v}$ bisects the angle $\measuredangle\left(v_{+}, A v_{+}\right)$, see $\operatorname{Tr}\left(v_{+}\right)=O T^{\prime} C^{\prime}$ on Figure 6 (a).

The middle vectors associated with a positive definite matrix play an essential role in the successful analysis of many matrix techniques used in Statistics and Numerical Analysis. Some aspects of this role are illustrated in Section 6.
4.6. Local optimisation when $A$ is indefinite. When $A$ is indefinite so is the sign of $g^{2}=$ $\lambda \lambda^{\prime}$. We restrict $A$ to be $2 \times 2: A_{0}=A_{\cdot \mathrm{M}}$ and $B_{0}=A_{0}-a I_{2}: \mathbf{M} \rightarrow \mathbf{M}$.

When $0<\lambda<\lambda^{\prime}$, the results of Corollary 4.9 apply readily to $A_{0}$. When $\lambda<\lambda^{\prime}<0,\left(-A_{0}\right)$ is positive definite and $\left(-A_{0}\right)^{1 / 2}$ is well-defined, $a$ and $h=\frac{g^{2}}{a}$ should be replaced by $|a|$ and $\frac{g^{2}}{|a|}$.

When $\lambda<0<\lambda^{\prime}, A_{0}$ is indefinite as well as $B_{0}$. However, the eigenvalues of $F_{0}=A_{0} B_{0}$ are positive, being $\left\{-e \lambda, e \lambda^{\prime}\right\}$.

The local analogue of Corollary 4.9 when $g^{2}<0$ is given by
Proposition 4.10. When $g^{2}<0$, let $\hat{v}$ be a middle vector. Then $v_{-}=|g| F_{0}^{-1 / 2} \hat{v}$ is an antieigenvector, $F_{0} v_{-}=|g| F_{0}^{1 / 2} \hat{v}$ and $\left\langle F_{0}^{-1 / 2} \hat{v}, F_{0}^{1 / 2} \hat{v}\right\rangle=-\frac{1}{g^{2}}\left\langle v_{-}, F_{0} v_{-}\right\rangle=1=\|\hat{v}\|^{2}$.

Proof. Straightforward calculation. See $\operatorname{Tr}\left(v_{-}\right)=O V^{\prime} C^{\prime}$ on Figure $6(\mathrm{~b})$ valid when $K=\mathbb{R}: \hat{v}$ bisects $\measuredangle\left(v_{-}, F_{0} v_{-}\right)$and $v_{-}, A v_{-}$are orthogonal $\left(\alpha=\frac{\pi}{2}\right)$.

(a) $g^{2}>0, a>0$
observation angle $\phi, \beta=\frac{\pi}{2}$

(b) $g^{2}<0, a>0$
perspective angle $\psi, \alpha=\frac{\pi}{2}$

Figure 6. $\operatorname{Tr}\left(v_{+}\right)(\mathrm{a})$ and $\operatorname{Tr}\left(v_{-}\right)(\mathrm{b})$ in $\mathbf{M}$ when $K=\mathbb{R}$

## 5. The 4D-invariant subspace M when $K=\mathbb{C}$

5.1. Angles between complex lines in real geometry. When $x$ and $y$ are nonzero complex vectors, the angle between the complex directions that they define is a subject that is rarely treated, even in advanced textbooks on linear algebra [Scharnhorst, 2001]. Since $\mathbb{C}^{n} \cong \mathbb{R}^{2 n}$ two complex lines define two real planes in a space of at least 4 real dimensions. Generically two Jordan (canonical) real angles are necessary to specify the relative position of two arbitrary real planes [Jordan, 1875]. In the present context where $A$ is hermitian, there is a vast geometric simplification. One can show that the two Jordan angles are equal [Kwietniewski, 1902, Maruyama, 1950, Wong, 1977, Theorem 1.7.4].
5.2. Angles between real planes in M. We consider the coupling $\lambda<\lambda^{\prime}$ and the associated 4D-subspace M spanned by the respective eigenvectors $q$ and $q^{\prime}$. Any $u=z q+z^{\prime} q^{\prime}$ belongs to the unit sphere in 4 dimensions $(S)=\left\{\left(z, z^{\prime}\right),|z|^{2}+\left|z^{\prime}\right|^{2}=1, z, z^{\prime} \in \mathbb{C}\right\}$. When $u$ is not an eigenvector and $a \neq 0$, the three complex vectors $a u, A u, B u$ define three real planes all passing through $O$. It is not an easy matter to interpret in 4 D the angles $\alpha, \beta$ and $\gamma$ which are Jordan canonical angles between the 3 planes. It is clear that the triangle $\operatorname{Tr}(u)$ which lies in the 2D-plane $\mathbf{M}$ when $K=\mathbb{R}$ and the trigonometric information it provides (Fig. 5 and 6 ) have no general counterparts when $K=\mathbb{C}$. Hence no known trigonometric interpretation is available to us in 4D. However, the properties of the triangle $\operatorname{Tr}(M)$ in the spectral plane cover both cases $K=\mathbb{R}$ and $\mathbb{C}$. Therefore a 2D-trigonometric interpretation involving ordinary angles remains available in the spectral plane $\boldsymbol{\Sigma}$, see Section 5.5.
5.3. The distinguished sets $D_{ \pm}, \hat{D}$ in $\mathbb{R}^{3}$. We recall that the cartesian product of two circles $S^{1} \times S^{1} \subset \mathbb{C}^{2}$ is homeomorphic, in topology, to a torus in $\mathbb{R}^{3}$. The three distinguished subsets $D_{ \pm}, \hat{D}$ of the unit sphere $(S)=S^{3} \subset \mathbb{R}^{4}$, which yield optimal angles in $\mathbf{M} \subset \mathbb{C}^{n}$, can therefore be interpreted as tori in $\mathbb{R}^{3}$. Indeed
i) $D_{+}=w_{+} S^{1} \times w_{+}^{\prime} S^{1}, w_{+}^{\prime}=\sqrt{1-w_{+}^{2}} \neq w_{+}, g^{2}>0$,
ii) $D_{-}=w_{-} S^{1} \times w_{-}^{\prime} S^{1}, w_{-}^{\prime}=\sqrt{1-w_{-}^{2}}, g^{2}<0$,
iii) $\hat{D}=\hat{w} S^{1} \times \hat{w} S^{1}, a \neq 0, \hat{w}=\frac{1}{\sqrt{2}}$.

In case i) or ii) with $a \neq 0$ the radii $w$ and $w^{\prime}$ are distinct and yield a ring torus. And in case ii) with $\lambda=-\lambda^{\prime}$ or iii) the equal radii $\frac{1}{\sqrt{2}}$ yield a horn torus, see Figure 7


Figure 7. Distinguished tori in $\mathbb{R}^{3}$ when $K=\mathbb{C}$
5.4. The locus of the end-point for $A u, u \in(S)$. Let $O \vec{M}^{\prime}=A u, u \in(S)$, we denote $M^{\prime}=\left(s, s^{\prime}\right)$ with $s, s^{\prime} \in \mathbb{C}$. Since $A u=\lambda z q+\lambda^{\prime} z^{\prime} q^{\prime}$ when $g^{2} \neq 0$, the point $M^{\prime}$ describes the quadratic surface $\frac{|s|^{2}}{\lambda^{2}}+\frac{\left|s^{\prime}\right|^{2}}{\lambda^{\prime 2}}=1$ in $\mathbb{R}^{4}$. If $g^{2}=0, \lambda=0<\lambda^{\prime}$ (say), $A u=\lambda^{\prime} z^{\prime} q^{\prime}$ and $M^{\prime}=\left(0, s^{\prime}\right)$ describes the disk $0 \times\left\{s^{\prime},\left|s^{\prime}\right| \leq \lambda^{\prime}\right\}$. Observe that the segment $\left[-\lambda^{\prime}, \lambda^{\prime}\right]$ when $K=\mathbb{R}$ becomes a disk when $K=\mathbb{C}$.

It follows readily that the endpoints of $A v_{+}$or $A v_{-}$describe the homothetic tori $g D_{+}$or $|g| D_{-}$for $g^{2} \neq 0$. If $g^{2}=0,\left\|A v_{+}\right\|=\left\|A v_{-}\right\|=0$ : the endpoints coalesce with the origin $O$.
5.5. The full trigonometric information is given by $\operatorname{Tr}(M)$. We show how the spectral plane is a complete source of information when $K=\mathbb{C}$. The angles, the distances and the weights describing the catchvectors (if $g^{2}>0$ ), the antieigenvectors (if $g^{2}<0$ ) or the middle vectors can be easily constructed by ruler and compass in the spectral plane. Given two eigenvalues $\lambda<\lambda^{\prime}$ and their relative position with respect to the origin $O$, we construct the circle $\Gamma$ (linking curve between $\lambda$ and $\lambda^{\prime}$ ) of radius $e$ and center $C$ with coordinates ( $a, 0$ ).


Figure 8. Sequential optimisation in the spectral plane, $g^{2} \neq 0$

The information corresponding to the catchvectors and antigeigenvectors contained in the spectral plane is given by the black triangles in Figure 8 (a) and (b) respectively. If $g^{2}>0$ (resp. $g^{2}<0$ ), we draw the tangent to the circle $\Gamma$ (resp. the vertical line) passing through $O$ which intersects $\Gamma$ at the point $T$ (resp. $V$ ) and we obtain simultaneously the angles $\alpha=$ $\angle(O C, O T)=\phi$ and $\beta=\angle(T C, T O)=\frac{\pi}{2}\left(\right.$ resp $\alpha=\angle(O C, O V)=\frac{\pi}{2}$ and $\beta=\angle(V C, V O)=$ $\psi$ ). Here we consider only the points $T$ (if $g^{2}>0$ ) and $V\left(\right.$ if $\left.g^{2}<0\right)$ but analogous results are obtained if we use the points $U$ and $W$ in Figure 2. Then, the angle $\gamma$ is obtained as follows: $\gamma=\angle(C O, C T)=\frac{\pi}{2}-\phi\left(\right.$ resp. $\left.\gamma=\angle(C O, C V)=\frac{\pi}{2}-\psi\right)$ if $g^{2}>0$ (resp. $g^{2}<0$ ), which is such that $\cos \gamma=\frac{e}{|a|}$ and $\sin \gamma=\frac{g}{|a|}$ (resp. $\cos \gamma=\frac{|a|}{e}$ and $\sin \gamma=\frac{|g|}{e}$ ). Finally, the angle $\delta$ is such that: $\delta=\frac{\gamma}{2}=\angle\left(\Lambda^{\prime} \Lambda, \Lambda^{\prime} M\right)$ (resp. $\delta=\frac{\gamma}{2}=\angle\left(\Lambda^{\prime} \Lambda, \Lambda^{\prime} V\right)$ ). Since the values $\cos \gamma$ and $\sin \gamma$ are known, by trigonometric identities we obtain that: $\cos \delta=\sqrt{\frac{\lambda^{\prime}}{2|a|}}=w_{+}$and $\sin \delta=\sqrt{\frac{\lambda}{2|a|}}=w_{+}^{\prime}$ (resp. $\cos \delta=\sqrt{\frac{\lambda^{\prime}}{2 e}}=w_{-}$and $\sin \delta=\sqrt{\frac{-\lambda}{2 e}}=w_{-}^{\prime}$ ) if $g^{2}>0$ (resp. $g^{2}<0$ ). Thus, the weights $w_{ \pm}, w_{ \pm}^{\prime}$ describing the sets $D_{ \pm}$in both the invariant plane $\mathbf{M}$ if $K=\mathbb{R}$ and the 4D-invariant subspace $\mathbf{M}$ if $K=\mathbb{C}$, are also found in the spectral plane.

In the case of middle vectors the information corresponds to the blue triangles with maximal surface $\frac{1}{2}|a| e$ in Figure 8 (a) and (b). We draw the vertical line passing through $C$ which intersects $\Gamma$ at $L^{+}=(a, e)$ and $L^{-}=(a,-e)$. In this case, $\gamma=\frac{\pi}{2}$ and $\delta=\frac{\pi}{4}$. Thus, the unique weight $\hat{w}=\frac{1}{\sqrt{2}}$ describing the set $\hat{D}$ is given by $\cos \delta=|\sin \delta|=\frac{1}{\sqrt{2}}$.

When $|a| \neq e \Leftrightarrow g^{2} \neq 0$, there exist two mutually exclusive optimal configurations: in one of them, $\alpha$ or $\beta=\frac{\pi}{2}$ according to the sign of $g^{2}$, in the other $\gamma=\frac{\pi}{2}$. This 2 -fold optimisation can be described as sequential, see Figure 8. When $|a|=e \Leftrightarrow g^{2}=0, \alpha=\beta$, it is possible to get the common value $\alpha=\beta=\gamma=\frac{\pi}{3}$ which is lesser than $\frac{\pi}{2}$, but is achieved by all three angles.

## 6. More on middle vectors when $A$ is positive definite

6.1. Numerical Analysis. As indicated in Remark 3.1, the consideration of the three angles $\alpha, \beta, \gamma$ sheds more light on the geometrical picture related to the well-known Wielandt and Kantorovich inequalities for a nonsingular matrix described in [Horn and Johnson, 1985, p. 441-445]. Here we consider the case where $A$ is a positive definite hermitian matrix with eigenvalues $0<\lambda_{1} \leq \ldots \leq \lambda_{n}$.

A relation between the maximal turning angle $\phi(A)$ and the Wielandt angle $\theta_{W}$ related to the condition number of a matrix has been given in [Gustafson, 1999]. The condition number of $A^{\frac{1}{2}}$ in the euclidean norm is $\operatorname{cond}\left(A^{\frac{1}{2}}\right)=\frac{\sqrt{\lambda_{n}}}{\sqrt{\lambda_{1}}}=\cot \angle\left(q, v_{+, *}\right)=\cot \theta_{+}$. The angle $\theta_{W}$ appearing in Wielandt's inequality is defined in the first quadrant by $\cot \left(\frac{\theta_{W}}{2}\right)=\operatorname{cond}\left(A^{\frac{1}{2}}\right)$ so that $\theta_{+}=\frac{\theta_{W}}{2}$. Proposition 4.8 in Section 4.4, tells us that $\alpha \leq \phi(A)$, hence $\theta_{W}=\frac{\pi}{2}-\phi(A)=\gamma(A)$. Thus, Wielandt's inequality for any pair of orthogonal vectors $x, y \in K^{n}$ and $A^{\frac{1}{2}}$ is given by

$$
\begin{equation*}
\frac{\left|\left\langle A^{\frac{1}{2}} x, A^{\frac{1}{2}} y\right\rangle\right|}{\left\|A^{\frac{1}{2}} x\right\|\left\|A^{\frac{1}{2}} y\right\|}=\frac{\langle x, A y\rangle}{\|x\|_{A}\|y\|_{A}} \leq \cos \theta_{W}=\cos \gamma(A)=\sin \phi(A) . \tag{6.1}
\end{equation*}
$$

The angles $\gamma(A)$ and $\phi(A)$ are complementary, thus if the turning angle $\phi(A)$ is large the angle $\theta_{W}=\gamma(A)$ is small which indicates that the matrix $A$ is ill-conditioned. The equality is attained if $x$ and $y$ are precisely orthogonal middle vectors in $\hat{D}_{*}$. The case $K=\mathbb{R}$ is illustrated in Figure 9 with $\hat{v}_{1}=\frac{1}{\sqrt{2}}\left(q_{n}+q_{1}\right)$ and $\hat{v}_{2}=\frac{1}{\sqrt{2}}\left(q_{n}-q_{1}\right)$, where $q_{1}$ and $q_{n}$ are eigenvectors associated with the eigenvalues $\lambda_{1}$ and $\lambda_{n}$ of $A$. Figure 9 shows the complementarity between $\gamma(A)$ and $\phi(A)$. The smaller the angle $\gamma(A)=\measuredangle\left(A^{\frac{1}{2}} \hat{v}_{1}, A^{\frac{1}{2}} \hat{v}_{2}\right)=\measuredangle\left(A v_{+, 1}, A v_{+, 2}\right)$, the closer the vectors $A^{\frac{1}{2}} \hat{v}_{1}$, and $A^{\frac{1}{2}} \hat{v}_{2}$ are to be parallel.

In [Horn and Johnson, 1985, p. 444], Kantorovich's inequality is derived from Wielandt's inequality (6.1). However, Kantorovich's inequality is equally a direct consequence of Proposition 4.8:

$$
\frac{\|x\|^{2}}{\|x\|_{A}\|x\|_{A^{-1}}} \geq \frac{2 \sqrt{\lambda_{1} \lambda_{n}}}{\lambda_{1}+\lambda_{n}}=\cos \phi(A)
$$

where equality is attained if $x$ is any middle vector $\hat{v}$ in $\hat{D}_{*}$. In Figure 9 the geometrical meaning of the angles $\phi(A)$ and $\theta_{W}=\gamma(A)$ when $K=\mathbb{R}$ is illustrated: $\frac{\|\hat{v}\|^{2}}{\|\hat{v}\|_{A}\| \|_{A^{-1}}}=$ $\cos \measuredangle\left(A^{-1 / 2} \hat{v}, A^{1 / 2} \hat{v}\right)=\cos \phi(A)$.
6.2. Statistics. Some of the matrix identities that are used in statistics and econometrics [Gustafson, 2002, 2012, chapter 6], [Wang and Chow, 1994] can benefit from the light provided by Section 4.5 . As a consequence of Corollary 4.9 we get

$$
\begin{equation*}
\cos \measuredangle\left(A^{-1 / 2} \hat{v}, A^{1 / 2} \hat{v}\right)=\cos \measuredangle\left(v_{+}, A v_{+}\right)=\frac{g}{a}, \tag{6.2}
\end{equation*}
$$

see Figure 9. It is shown in [Rincon-Camacho, 2015a, b] that (6.2) provides the missing step to elucidate the worst case of statistical efficiency presented in [Gustafson, 2002, pp. 147150], [Gustafson, 2012, theorem 6.2, pp. 102-103]. The equality (6.2) tells us that the worst inefficiency is characterised equally by the catchvectors $v_{+}$by means of $\phi=\measuredangle\left(v_{+}, A_{+}\right)$or by the middle vectors $\hat{v}$ through $\measuredangle\left(A^{-1 / 2} \hat{v}, A^{1 / 2} \hat{v}\right)$. These middle vectors have been introduced
as "inefficient" vectors in [Gustafson, 2002, 2012]. For a detailed analysis of this important application to statistics, see [Rincon-Camacho, 2015a].


Figure 9. A middle vector pair in $\hat{D}_{*}, K=\mathbb{R}$

## 7. A TRIGONOMETRIC ASPECT FOR $K=\mathbb{C}^{n}$

7.1. The source of Gustafson's theory. Assuming that $A$ is symmetric positive definite in $\mathbb{R}^{n \times n}$, Gustafson added to the functional:

$$
x \in \mathbb{R}^{n},\|x\|=1 \mapsto c(x)=\frac{x^{T} A x}{\|A x\|} \in \mathbb{R}^{+}
$$

the companion functional:

$$
(\eta, x) \in \mathbb{R}^{+} \times \mathbb{R}^{n},\|x\|=1 \mapsto n(\eta, x)=\|(\eta A-I) x\| \in \mathbb{R}^{+} .
$$

This led him to the remarkable theorem

$$
\begin{equation*}
\max _{\|x\|=1} \min _{\eta>0} n(\eta, x)=\min _{\eta>0} \max _{\|x\|=1} n(\eta, x)=\frac{\lambda_{n}-\lambda_{1}}{\lambda_{1}+\lambda_{n}}<1, \tag{7.1}
\end{equation*}
$$

where $0<\lambda_{1}<\lambda_{n}$ are the extreme eigenvalues of $A$. Since $\cos \phi(A)=\frac{2 \sqrt{\lambda_{1} \lambda_{n}}}{\lambda_{1}+\lambda_{n}}=\frac{g_{*}}{a_{*}}$, it becomes clear that the value of (7.1) is but $\sin \phi(A)=\frac{e_{*}}{a_{*}}$ where $\phi(A)$ denotes the maximal turning angle for $A$ [Gustafson, 1968]. Actually (7.1) is a by-product of Gustafson's much more encompassing work on perturbation theory of semi-groups, see Chapter 1 in [Gustafson, 2012]. And Gustafson insists repeatedly that one has no trigonometry unless one has a cosine and a sine.

Our purpose in this Section is to examine how we can get a trigonometry in the general case $A$ hermitian indefinite over $K=\mathbb{C}$. This we do by first considering the functional:

$$
\begin{equation*}
(\eta, x) \in \mathbb{R} \times \mathbb{C}^{n},\|x\|=1 \mapsto \nu(\eta, x)=\|(\eta Y Z-I) x\|^{2} \tag{7.2}
\end{equation*}
$$

where $Y$ and $Z$ are two commuting hermitian matrices, $S=Y Z=Z Y$.
7.2. The partial derivative $\frac{\partial \nu}{\partial \eta}(\eta, x)=0,\|x\|=1$. We write

$$
\nu(\eta, x)=\langle\eta Y Z x-x, \eta Y Z x-x\rangle=\|Y Z x\|^{2} \eta^{2}-2\langle x, Y Z x\rangle \eta+1, \text { for }\|x\|=1
$$

Then $\frac{\partial \nu}{\partial \eta}=2\left(\|Y Z x\|^{2} \eta-\langle x, Y Z x\rangle\right)=0$ iff $\eta$ equals $\eta_{0}(x)=\frac{\langle x, Y Z x\rangle}{\|Y Z x\|^{2}}$ defined for $x$ such that $Y Z x \neq 0$. We set $\jmath(x)=\frac{\langle Y x, Z x\rangle}{\|Y Z x\|},|\jmath(x)| \in[0,1]$.

Lemma 7.1. For $x$ fixed in $K^{n},\|x\|=1$, the condition $\frac{\partial \nu}{\partial \eta}=0$ at $\left(\eta_{0}, x\right)$ entails that $\nu_{0}(x)=$ $\nu\left(\eta_{0}, x\right)=1-\jmath^{2}(x), 0 \leq \nu_{0}(x) \leq 1$.

Proof. Clear: set $\eta_{0}=\frac{\partial(x)}{\|Y Z x\|}$ in $\nu(\eta, x)$. We get $\nu_{0}(x)=1-\left(\frac{x^{H} Y Z x}{\|Y Z x\|}\right)^{2},\|x\|=1, Y Z x \neq 0$, that is $1-\jmath^{2}(x)$.

When can we relate $\jmath(x)$, the cosine of angle $(x, S x)$ to that of one of the three angles of interest $\alpha, \beta$ or $\gamma$ in $\boldsymbol{\Sigma}$ ? In other words, when is it that $\|Y Z x\|=\|Y x\|\|Z x\|$ ?

When $Y$ is proportional to $I$, the answer is obvious:

- $Y=I, Z=A$ yield $\jmath(x)=\cos \alpha$,
- $Y=-I, Z=B$ yield $\jmath(x)=\cos \gamma$.

When $Y=B, Z=A, F=B A$, the answer is given by the
Lemma 7.2. For any $u \in(S) \subset \mathbf{M}, \jmath(u)=\cos \beta$.
Proof. We have to show that $\|B A u\|=\|B u\|\|A u\|$ for $u \in(S)$. With $A u=e^{i \xi} \lambda q+e^{i \xi^{\prime}} \lambda^{\prime} q^{\prime}$, we get $B A u=e\left(-e^{i \xi} \lambda q+e^{i \xi^{\prime}} \lambda^{\prime} q^{\prime}\right)$, hence $\|B A u\|=e\|A u\|$, $e=\|B u\|$.
7.3. The trigonometric aspect $\zeta:\|x\|=1 \mapsto \zeta(x) \in\left[0, \frac{\pi}{2}\right]$. Let $\mathbf{S}$ be the unit sphere in $\mathbb{C}^{n}$ we consider the two connected functions $\mathbf{S} \rightarrow[0,1]$ defined respectively by $\kappa: x \in \mathbf{S} \mapsto \kappa(x)=$ $|\jmath(x)|$ and $\sigma: x \in \mathbf{S} \mapsto \sigma(x)=\sqrt{\nu_{0}(x)}$. The two functions are connected by the relation

$$
\begin{equation*}
\kappa^{2}(x)+\sigma^{2}(x)=1 \text { for all } x \text { in } \mathbf{S} \tag{7.3}
\end{equation*}
$$

In the plane $\mathbb{R}^{2}$, the numbers $\kappa$ and $\sigma$ in $[0,1]$ define $\zeta$ in $\left[0, \frac{\pi}{2}\right]$ such that $\kappa=\cos \zeta$ and $\sigma=\sin \zeta$. See Figure 10.

Definition 7.1. The function $\zeta: \mathbf{S} \rightarrow\left[0, \frac{\pi}{2}\right]$ is called aspect of $x \in \mathbb{C}^{n}$. It is defined by the identities $\kappa(x)=\cos \zeta(x)=|\jmath(x)|$ and $\sigma(x)=\sin \zeta(x)=\sqrt{\nu_{0}(x)}$.

The concept of aspect is an emerging one: it captures the essence of the identity (7.3) which links $\kappa(x)$ and $\sigma(x): \zeta(x)$ is the ordinary angle in the quarter plane defined by points on the unit circle with nonnegative components $(\kappa(x), \sigma(x))$, see Figure 10. It provides the geometric meaning of an angle in $\mathbb{R}^{2}$ to the analytic variable $\zeta(x)=\operatorname{angle}(x, S x) \in\left[0, \frac{\pi}{2}\right]$ when $x \in \mathbb{C}^{n}$. By comparison, it is an easy matter, when $x \in \mathbb{R}^{n}$ to relate $\zeta(x)$ to $\mathcal{S}(x)=\measuredangle(x, S x) \in[0,2 \pi]$.


Figure 10. The aspect $\zeta$ of $x \in \mathbb{C}^{n}$
7.4. Three applications of Lemma 7.1. We address the three possibilities that $u \in D_{ \pm}, \hat{D}$ in $\mathbf{M} \subset \mathbb{C}^{n}$.

Proposition 7.3. The distinguished sets yield the following results where $u \in(S)$ :
i) $v_{+} \in D_{+} \Rightarrow \eta_{0}\left(v_{+}\right)=\frac{1}{a}, \nu_{0}(u)=\frac{e^{2}}{a^{2}}=\sin ^{2} \phi$,
ii) $v_{-} \in D_{-} \Rightarrow \eta_{0}\left(v_{-}\right)=\frac{1}{e^{2}}>0, \nu_{0}(u)=\frac{a^{2}}{e^{2}}=\sin ^{2} \psi$,
iii) $\hat{v} \in \hat{D} \Rightarrow \eta_{0}(\hat{v})=0, \nu_{0}(u)=1$.

Proof. i) $\eta_{0}\left(v_{+}\right)=\frac{1}{a}$ follows from $v_{+}^{H} A v_{+}=\frac{g^{2}}{a}$ and $\left\|A v_{+}\right\|=g ; \nu_{0}(u)=\left\|\left(\frac{1}{a} A-I\right) u\right\|^{2}=$ $\frac{1}{a^{2}}\|B u\|^{2}$.
ii) $\eta_{0}\left(v_{-}\right)=\frac{1}{e^{2}}$ follows from $v_{-}^{H} F v_{-}=-g^{2}>0$ and $\left\|F v_{-}\right\|^{2}=-g^{2} e^{2} ; \nu_{0}(u)=\left\|\left(\frac{1}{e^{2}} F-I\right) u\right\|^{2}=$ $\frac{a^{2} e^{2}}{e^{4}}=\frac{a^{2}}{e^{2}}$.
iii) $\eta_{0}(\hat{v})=0$ follows from $\langle\hat{v}, B \hat{v}\rangle=0 ; \nu_{0}(u)=\|-u\|=1$.

It is clear that even though $\nu_{0}(u)$ is constant at the optimal value for any $u \in(S)$, the aspect $\zeta$ emerges only for vectors in the corresponding distinguished set. When this occurs, $\zeta$ can be identified with any of the angles $\alpha, \beta, \gamma$ in $\boldsymbol{\Sigma}$ provided one of them is $\frac{\pi}{2}$.

To proceed toward a min-max equality, it is useful to introduce a $2 \times 2$ hermitian matrix $H$ with eigenvalues $h<h^{\prime}$, and to consider $n(\eta, x)=\sqrt{\nu(\eta, x)}$

Lemma 7.4. i) If $h h^{\prime}>0$, then

$$
\max _{x \in \mathbb{C}^{2},\|x\|=1} \min _{\eta \in \mathbb{R}}\left\|\left(\eta H-I_{2}\right) x\right\|=\min _{\eta}\left\|\eta H-I_{2}\right\|=\frac{h^{\prime}-h}{\left|h^{\prime}+h\right|} .
$$

ii) If $h<0<h^{\prime}$ and $h+h^{\prime}=0$, then

$$
\max _{x \in \mathbb{C}^{2},\|x\|=1} \min _{\eta \in \mathbb{R}}\left\|\left(\eta H-I_{2}\right) x\right\|=\min _{\eta}\left\|\eta H-I_{2}\right\|=1 .
$$

Proof. i) $\left\|\eta H-I_{2}\right\|=\max \left(|\eta h-1|,\left|\eta h^{\prime}-1\right|\right)$. The minimal value for $\eta \in \mathbb{R}$ is achieved when the two quantities to be compared are equal, hence $\eta_{0}=\frac{2}{h+h^{\prime}}$, yielding the right-hand side equality at $\frac{h^{\prime}-h}{\left|h^{\prime}+h\right|}>0$.

We turn to $\min _{\eta}\left\|\left(\eta H-I_{2}\right) x\right\|$ with $H=Q\left(\begin{array}{cc}h & 0 \\ 0 & h^{\prime}\end{array}\right) Q^{H},\|x\|=1$. We set $\Delta_{0}=\operatorname{diag}\left(h, h^{\prime}\right)$, $n(\eta, x)=\left\|\left(\eta H-I_{2}\right) x\right\|=\left\|\left(\eta A_{0}-I_{2}\right) Q^{H} x\right\|, y=Q^{H} x=\left(y_{1}, y_{2}\right)^{T} \in \mathbb{C}^{2}, Y(\eta, y)=((\eta h-$ 1) $\left.y_{1},\left(\eta h^{\prime}-1\right) y_{2}\right)^{T}$ and $n^{2}(\eta, x)=N(\eta, y)=\|Y(\eta, y)\|^{2}=(\eta h-1)^{2} y_{1}^{2}+\left(\eta h^{\prime}-1\right)^{2} y_{2}^{2}$. Thus $N(\eta, y)=\mathbf{A} \eta^{2}-2 \mathbf{B} \eta+1$ with $\mathbf{A}(y)=h^{2} y_{1}^{2}+h^{\prime 2} y_{2}^{2}=\left\|\Delta_{0} y\right\|^{2}>0, \mathbf{B}(y)=h y_{1}^{2}+h^{\prime} y_{2}^{2}=\left\langle y, \Delta_{0} y\right\rangle$. It follows that $N(\eta, y) \geq 0$ for $\eta \in \mathbb{R}, y \in \mathbb{C}^{2}$. The 0 value is achieved if $\mathbf{B}^{2}=\mathbf{A} \Leftrightarrow x$ is an eigenvector. When $x$ is not an eigenvector, $N(\eta, y)$ is positive and minimum for $y$ fixed if $\eta$ solves $\frac{\partial N}{\partial \eta}=0$, that is $\eta_{1}(y)=\frac{\mathbf{B}}{\mathbf{A}}$. Hence $\min _{\eta} N(\eta, y)=N\left(\eta_{1}, y\right)=N_{1}(y)=1-\frac{\mathbf{B}^{2}}{\mathbf{A}}=$ $n_{1}^{2}(x)=\min _{\eta}\left\|\left(\eta H-I_{2}\right) x\right\|^{2}$. Therefore $\left(\max _{\|y\|=1}\left(1-\frac{\mathbf{B}^{2}}{\mathbf{A}}\right)<1\right)=1-\min _{\|y\|=1}\left(\frac{\left\langle y, \Delta_{0} y\right\rangle}{\left\|\Delta_{0} y\right\|}\right)^{2}=$ $1-\frac{4 h h^{\prime}}{\left(h^{\prime}+h\right)^{2}}=\left(\frac{h^{\prime}-h}{h^{\prime}+h}\right)^{2}$ is achieved for $y_{+}=\left(e^{i \xi} w_{+}, e^{i \xi^{\prime}} w_{+}^{\prime}\right)^{T}=Q^{H} v_{+}, v_{+} \in D_{+}$. Note that

$$
\eta_{1}\left(y_{+}\right)=\frac{\left\langle y_{+}, \Delta_{0} y_{+}\right\rangle}{\left\|\Delta_{0} y_{+}\right\|^{2}}=\frac{\left\langle v_{+}, H v_{+}\right\rangle}{\left\|H v_{+}\right\|^{2}}=\frac{2 h h^{\prime}}{\left(h+h^{\prime}\right) h h^{\prime}}=\eta_{0} .
$$

ii) $\left\|\eta H-I_{2}\right\|=\max \left(\left|\eta h^{\prime}-1\right|,\left|\eta h^{\prime}+1\right|\right) \geq 1$ since $h^{\prime}=-h>0$. The minimal value 1 is achieved for $\eta=0$. On the other hand, let $x \in \hat{D}$, then $\left|y_{1}\right|=\left|y_{2}\right|=\frac{1}{\sqrt{2}}$ and $\mathbf{B}=\left\langle y, \Delta_{0} y\right\rangle=$ $\frac{1}{\sqrt{2}}\left(h+h^{\prime}\right)=0$. Thus $1-\frac{\mathbf{B}^{2}}{\mathbf{A}}=1$.

Example 7.1. We illustrate Lemma 7.4 i) on $H=\left(\begin{array}{cc}0.5 & -0.3 \\ -0.3 & 0.5\end{array}\right)$ with eigenvalues $0<h=$ $0.2<h^{\prime}=0.8$ and $Q=\frac{1}{\sqrt{2}}\left(\begin{array}{cc}1 & 1 \\ 1 & -1\end{array}\right)$. For $x=\cos \theta q+\sin \theta q^{\prime}, y=Q^{T} x=(\cos \theta, \sin \theta)^{T}$ : the point $(\eta, x)$ is coded in $\mathbb{R}^{2}$ as $\eta y=(\eta \cos \theta, \eta \sin \theta)$. See on Figure 11 (a) a sample of 5 circles with increasing radius $\eta$. With $a=\frac{1}{2}, e=0.3, w_{+}=\sqrt{0.8} \sim 0.89, w_{+}^{\prime}=\sqrt{0.2} \sim 0.45$, $\eta_{1}\left(y_{+}\right)=\eta_{0}=2, n(2, x)=\frac{0.3}{0.5}=0.6=\sin \phi(H)$ where $\phi(H)$ is the turning angle, finally $\tan \theta_{+}=\frac{w_{+}^{\prime}}{w_{+}}=\sqrt{\frac{h}{h^{\prime}}}=\sqrt{0.25}=0.5$. See Figure 11 (b).


Figure 11. (a) Code for $(\eta, x) \cong(n y) \in \mathbb{R}^{2}$ (b) triangle $\operatorname{Tr}\left(v_{+}\right)$in $\mathbb{R}^{2}$
Figure 12 displays for $\eta>0$ the corresponding functional $n(\eta, x)=\left\|\left(\eta H-I_{2}\right) x\right\|, x \in \mathbb{R}^{2}$, $\|x\|=1$ for $0 \leq \eta \leq 4$ and $\theta \in[0,2 \pi]$ in three views: (a) (resp. (b), (c)) displays the surface $(\eta y, n(\eta, x))$ in 3D (resp. (b) the 2D view from above, (c) the side view along $\theta_{+} \sim 0.46365$ ).

The saddle points correspond to $\eta_{1}\left(y_{+}\right)=\eta_{0}$ : they are the intersection of the critical value $\eta_{1}$ (in black) and the angle $\theta_{+}$(in red) which defines the 4 catchvectors $v_{+}=\varepsilon \cos \theta_{+} q+\varepsilon^{\prime} \sin \theta_{+} q^{\prime}$, $\varepsilon, \varepsilon^{\prime}= \pm 1$ in $D_{+}$, see Figure 12 (a) and (b). We observe that for $\eta_{1}$ fixed at $\eta_{1}(y)$, the value $n\left(\eta_{1}(y), y\right)=\sqrt{1-\left(\frac{\left\langle y, \Delta_{0} y\right\rangle}{\left\|\Delta_{0} y\right\|}\right)^{2}}=\sin \angle\left(y, \frac{\Delta_{0} y}{\left\|\Delta_{0} y\right\|}\right) \leq \sin \phi(H)=0.6$ and for $\theta_{+}$fixed, $n\left(\eta, v_{+}\right)$ describes the piecewise convex curve depending on $\eta$ with minimum value at $\eta_{1}\left(y_{+}\right)=\eta_{0}=2$, displayed in red on Figure 12 (c).

We recall that $A_{0}=A_{\lceil\mathbf{M}}, B_{0}=A_{0}-a I_{2}, F_{0}=A_{0} B_{0}$.
Corollary 7.5. Let be given a coupling $\lambda<\lambda^{\prime}, g^{2}=\lambda \lambda^{\prime}$. For $u \in(S) \subset \mathbf{M}$, the following min-max equalities hold:
i) if $g^{2}>0 \Leftrightarrow|a|>e,(7.4) \Leftrightarrow$
$\max _{u} \min _{\eta}\|(\eta A-I) u\|=\min _{\eta} \max _{u}\|(\eta A-I) u\|=\min _{\eta}\left\|(\eta A-I)_{\mid \mathbf{M}}\right\|=\left\|\frac{1}{a} A_{0}-I_{2}\right\|=\frac{e}{a}=\sin \phi$,
ii) if $g^{2}<0 \Leftrightarrow|a|<e,(7.5) \Leftrightarrow$ $\max _{u} \min _{\eta}\|(\eta F-I) u\|=\min _{\eta} \max _{u}\|(\eta F-I) u\|=\min _{\eta}\left\|(\eta F-I)_{\mid \mathbf{M}}\right\|=\left\|\frac{1}{e^{2}} F_{0}-I_{2}\right\|=\frac{|a|}{e}=\sin \psi$, iii) if $g^{2} \neq 0,(7.6) \Leftrightarrow$

$$
\max _{u} \min _{\eta}\|(\eta B-I) u\|=\min _{\eta} \max _{u}\|(\eta B-I) u\|=\min _{\eta}\left\|(\eta B-I)_{\mid \mathbf{M}}\right\|=\min _{\eta}\left\|\eta B_{0}-I_{2}\right\|
$$

$$
=\left\|I_{2}\right\|=1=\sin \frac{\pi}{2} .
$$

Proof. (7.4) and (7.5) follow from Lemma 7.4 i) since $\lambda \lambda^{\prime}>0$ for $A_{0}$ and $-e \lambda \lambda^{\prime}>0$ for $F_{0}$ (with eigenvalues $\left\{-e \lambda, e \lambda^{\prime}\right\}$ ). And (7.6) follows from Lemma 7.4 ii) since the eigenvalues of $B_{0}$ are $\{-e, e\}$ with zero mean: $-e+e=0$.


Figure 12. $n(\eta, x)=\left\|\left(\eta H-I_{2}\right) x\right\|$ (a) Surface (b) Top view (c) Side view $\theta_{+}$

When $A$ is positive definite (7.1) is a direct consequence of (7.4).
The following general remark is in order. Given the coupling $\left(\lambda, \lambda^{\prime}\right)$ characterised by ( $a=$ $\frac{\lambda+\lambda^{\prime}}{2}, e=\frac{\lambda^{\prime}-\lambda}{2}$ ), the maximal value $\sin \phi$ if $g^{2}>0$ (resp. $\sin \psi$ if $g^{2}<0$ ) is the ratio $\frac{e}{|a|}$ (resp. $\frac{|a|}{e}$ ) related to a 2 or 4 D -optimisation process in $\mathbf{M} \subset K^{n}$, whereas the maximal surface of the rectangle with diagonal $O L^{+}=\sqrt{a^{2}+e^{2}}$ is the product $|a| e$. Moreover, the ratios can be equally related to a 3 or 5 D min-max optimisation process taking place in $\mathbb{R} \times \mathbf{M}, \mathbf{M} \subset K^{n}$, leading to constant values in M. Note that $g^{2}=0 \Leftrightarrow \frac{e}{|a|}=\frac{|a|}{e}=1$ and $|a| e=e^{2}$.

## 8. Conclusion and perspective on spectral information processing

8.1. Summary. We have seen that the simultaneous consideration of the eigenelements $(\lambda, q)$ and $\left(\lambda^{\prime}, q^{\prime}\right)$ for hermitian matrices opens new vistas about their spectral theory, a domain considered so-far as almost completely researched. Spectral information is processed in a way which mixes analysis, elementary plane geometry and trigonometry in a most elegant and rich fashion. The added value is that the proofs are embarrassingly simple! If $A$ invertible admits $d$ distinct eigenvalues, $2 \leq d \leq n$, there are $d \frac{d-1}{2}$ eigenvalue pairs $\lambda<\lambda^{\prime}$ which produce $d(d-1)$ new informations in the form of the ratios $\frac{e}{|a|}\left(g^{2}>0\right)$ or $\frac{|a|}{e}\left(g^{2}<0\right)$ and the products $|a| e$. Most interestingly, the information process resulting from spectral coupling is delivered under various geometric forms, metric and trigonometric, derived from various right-angled triangles. To the matrix $A$, each coupling implicitly adds the two auxiliary matrices $B=A-a I$ and $F=A B=B A$ providing optimal orthogonality: $\left\langle v_{+}, F v_{+}\right\rangle=0$ in $D_{+}$if $g^{2}>0,\left\langle v_{-}, A v_{-}\right\rangle=0$ in $D_{-}$if $g^{2}<0$ and $\langle\hat{v}, B \hat{v}\rangle=0$ in $\hat{D}$. Therefore six (resp. four) distinguished aspects are associated with any given coupling such that $g^{2} \neq 0$ (resp. $=0$ ), according to the table

| $g^{2}$ | - | + | 0 |
| :---: | :---: | :---: | :---: |
| $\alpha, \gamma, \gamma$ | $\frac{\pi}{2}, \psi, \frac{\pi}{2}-\psi$ | $\phi, \frac{\pi}{2}, \frac{\pi}{2}-\phi$ | $\frac{\pi}{2}, \frac{\pi}{2}, 0$ |
|  | and | and |  |
|  | $\hat{\alpha}, \hat{\beta}, \frac{\pi}{2}$ | $\frac{\pi}{4}, \frac{\pi}{4}, \frac{\pi}{2}$ |  |

Why is it that the visionary insight of Gustafson - dating back from almost half a centuryhas attracted so little attention? A possible answer is that it explains more elegantly a body of well-known techniques which had already been proved to work by more familiar arguments -both in theory and practice -. The work is regarded by many as a nice confirmation of established facts. Only a handful of scientists have been receptive to its implicit potential for advancing our knowledge about the natural information processing realised through hermitian matrices.
8.2. Spectral information processing in the bireal plane. It has been proved elsewhere [Chatelin, 2016] that the way the information is optimally processed, either in the spectral plane or in the invariant plane if $K=\mathbb{R}$ is the result of an underlying multiplication in $\mathbb{R}^{2}$, known as bireal (Cockle 1848) which differs from the complex one (Cardano 1545). In $\mathbb{C}$, the non real unit is the imaginary number $i=\sqrt{-1}, i^{2}=-1$. By contrast, in the ring ${ }^{2} \mathbb{R}$ of bireal numbers, the non real unit is the unipotent number $\mathbf{u}, \mathbf{u}^{2}=1(\mathbf{u} \neq \pm 1)$. Thus if $z=x+\mathbf{u} y, z \times z^{\prime}=(x+y \mathbf{u}) \times\left(x^{\prime}+y^{\prime} \mathbf{u}\right)=x x^{\prime}+y y^{\prime}+\left(x^{\prime} y+x y^{\prime}\right) \mathbf{u}$, and $z^{*}=x-y \mathbf{u}$ [Rincon-Camacho and Latre, 2013]. Note that the above use of the term "unipotent" departs from the conventional one in abstract algebra where $(\mathbf{u}-1)^{n}=0$. The indefinite quadratic form $z \times z^{*}=x^{2}-y^{2} \in \mathbb{R}$ is nonzero for $|x| \neq|y|$ only: the algebra ${ }^{2} \mathbb{R}=\mathbb{R} \oplus \mathbf{u} \mathbb{R}$ has a ring structure with zerodivisors equipped with the hyperbolic measure $\sqrt{\left|x^{2}-y^{2}\right|}$ which is not a norm in the mathematical sense. By comparison $\mathbb{C}=\mathbb{R} \oplus i \mathbb{R}$ has a field structure equipped with the modulus $\sqrt{x^{2}+y^{2}}$ which is the true euclidean norm $|x+i y|=\|z\|$. This may be the reason why $\mathbb{C}$ is the preferred structure to equip $\mathbb{R}^{2}$ in modern Calculus. But this choice narrows dramatically our understanding of natural information processing. Such a narrowing is amply documented in [Chatelin, 2016], under its two-fold aspect, algebraic and analytic.
8.3. Bireal roots of the characteristic polynomial. Because the $n$ eigenvalues of $A$ hermitian are real its characteristic polynomial has real coefficients. As such it can be regarded as an element in $\mathbb{C}[X]$ as well as in ${ }^{2} \mathbb{R}[X]$. We know that the first choice provides no additional information to the real spectrum. But the second choice does provide new bireal eigenvalues! This follows readily from

$$
\begin{aligned}
0 & =\mu^{2}-2 a \mu+g^{2}=(\mu-\lambda)\left(\mu-\lambda^{\prime}\right) \text { for } \mu \in \mathbb{R} \text { or } \mathbb{C} \\
& =(\mu-(a+e \mathbf{u})) \times(\mu-(a-e \mathbf{u})) \text { for } \mu \in^{2} \mathbb{R}
\end{aligned}
$$

In other words the real eigenpair $\left(\lambda, \lambda^{\prime}\right)$ is complemented by the bireal eigenpair ( $\sigma=a+e \mathbf{u}$, $\sigma^{*}=a-e \mathbf{u}$ ) which are represented on Figures 2 and 8 by the points $L^{+}$and $L^{-}$of $\Gamma$, if one agrees to equip the spectral plane $\boldsymbol{\Sigma}=\mathbb{R}^{2}$ with the bireal ring structure ${ }^{2} \mathbb{R}$.

Of course, it is not surprising that a polynomial of degree $n$ may have more than $n$ roots in a ring, rather than the $n$ roots classically expected in an algebraically closed field such as $\mathbb{C} \supset \mathbb{R}$.

When $K=\mathbb{R}$ there is more to the bireal story: the invariant plane $\mathbf{M}$ itself can be thought of algebraically as the ring ${ }^{2} \mathbb{R}$. If one sets $q \cong 1, q^{\prime} \cong \mathbf{u}$, then the two middle vectors $\hat{v}_{1}=\frac{1}{\sqrt{2}}\left(q+q^{\prime}\right)=\sqrt{2} \frac{1+\mathbf{u}}{2}$ and $\hat{v}_{2}=\frac{1}{\sqrt{2}}\left(q-q^{\prime}\right)=\sqrt{2} \frac{1-\mathbf{u}}{2}$ are zerodivisors: $\hat{v}_{1} \times \hat{v}_{2}=\frac{1}{2}\left(1-\mathbf{u}^{2}\right)=0$. They are proportional to the idempotent numbers $\mathbf{e}_{ \pm}=\frac{1}{2}(1 \pm \mathbf{u})$ which satisfy $\mathbf{e}_{ \pm}^{2}=\mathbf{e}_{ \pm}$; thus $\hat{v}^{2}=\sqrt{2} \hat{v}$ for $\hat{v} \in \hat{D}$. In the idempotent basis $\left\{\mathbf{e}_{+}, \mathbf{e}_{-}\right\}$the bireal multiplication is performed componentwise. Let $z=x+y \mathbf{u}=\frac{x+y}{2} \mathbf{e}_{+}+\frac{x-y}{2} \mathbf{e}_{-}=X \mathbf{e}_{+}+Y \mathbf{e}_{-}$, then $z \times z^{\prime}=$ $\left(X \mathbf{e}_{+}+Y \mathbf{e}_{-}\right) \times\left(X^{\prime} \mathbf{e}_{+}+Y^{\prime} \mathbf{e}_{-}\right)=X X^{\prime} \mathbf{e}_{+}+Y Y^{\prime} \mathbf{e}_{-}$.

However the idempotent basis is orthogonal but not normalised: $\left\|\mathbf{e}_{ \pm}\right\|=\frac{1}{\sqrt{2}}$. The two middle vectors $\hat{v}_{1}=\sqrt{2} \mathbf{e}_{+}, \hat{v}_{2}=\sqrt{2} \mathbf{e}_{-}$provide an orthonormal basis. If $z=X \mathbf{e}_{+}+Y \mathbf{e}_{-}=\frac{1}{\sqrt{2}}\left(X \hat{v}_{1}+\right.$ $\left.Y \hat{v}_{2}\right)=\hat{X} \hat{v}_{1}+\hat{Y} \hat{v}_{2}$, the multiplication $\times$ is written in $\left\{\hat{v}_{1}, \hat{v}_{2}\right\}$ as $z \times z^{\prime}=\sqrt{2}\left(\hat{X} \hat{X}^{\prime} \hat{v}_{1}+\hat{Y} \hat{Y}^{\prime} \hat{v}_{2}\right)$.
"And what if $K=\mathbb{C}$ ?" will rightly ask the curious reader. Just like ${ }^{2} \mathbb{R}$ can replace $\mathbb{C}$ for multiplication in $\mathbb{R}^{2}$, in the four real dimensions of $\mathbb{R}^{4} \cong \mathbb{C}^{2}$, the classical noncommutative field
$\mathbb{H}$ of quaternions can be replaced by the commutative ring of bicomplex numbers ${ }^{2} \mathbb{C}=\mathbb{C} \oplus \mathbf{u} \mathbb{C}$, where $\mathbf{u}$ is now a 4 D -unipotent vector $\left(\mathbf{u}^{2}=1, \mathbf{u} \neq \pm 1\right)$. For more on bicomplex multiplication, see [Cockle, 1848, Segre, 1892, Price, 1991, Rincon-Camacho and Latre, 2013].
8.4. Bireals for hydrodynamics. Some physicists with a strong mathematical leaning are aware of the computational potential that resides in the inconspicuous ring of 2D-bireal numbers. The book [Lavrentiev and Chabat, 1980], a classic in hydrodynamics, devotes its whole chapter 2 more generally to the three families of algebraic structures that can equip the real plane. It also compares the properties of functions in 2 real variables which are analytic with respect to the complex or the bireal multiplication, the latter being useful to model supersonic flows, see chapter 4 of the same book. It is possible that the information carried by spectral coupling could bring new insights into computational fluid dynamics. The matter is investigated elsewhere [Chatelin, 2016].

## References

[Chaitin-Chatelin and Frayssé, 1996] Chaitin-Chatelin, F. and Frayssé, V. (1996). Lectures on finite precision computations. Software, Environments, Tools, volume 1. SIAM.
[Chatelin, 2012a] Chatelin, F. (2012a). Eigenvalues of Matrices: Revised Edition. Classics in Applied mathematics, volume 71. SIAM.
[Chatelin, 2012b] Chatelin, F. (2012b). Qualitative Computing: a computational journey into nonlinearity. World Scientific, Singapore.
[Chatelin, 2016] Chatelin, F. (2016). Numbers in Mind: from Information to Meaning. An essay on the transformative dynamics of multiplication (in preparation). World Scientific and Imperial College Press, London.
[Cockle, 1848] Cockle, J. (1848). III. On Certain Functions Resembling Quaternions, and on a New Imaginary Algebra. Phil. Mag. (3), 33:435-439.
[Gustafson, 1968] Gustafson, K. (1968). A min-max theorem. Notices Amer. Math. Soc, 15:799.
[Gustafson, 1999] Gustafson, K. (1999). The geometrical meaning of the Kantorovich-Wielandt inequalities. Linear Algebra and its Applications, 296(1):143-151.
[Gustafson, 2002] Gustafson, K. (2002). Operator trigonometry of statistics and econometrics. Linear Algebra and its Applications, 354(1):141-158.
[Gustafson, 2012] Gustafson, K. (2012). Antieigenvalue Analysis: With Applications to Numerical Analysis, Wavelets, Statistics, Quantum Mechanics, Finance and Optimization. World Scientific.
[Horn and Johnson, 1985] Horn, R. A. and Johnson, C. R. (1985). Matrix analysis. Cambridge University Press.
[Jordan, 1875] Jordan, C. (1875). Essai sur la géométrie à $n$ dimensions. Bulletin de la Société mathématique de France, 3:103-174.
[Kwietniewski, 1902] Kwietniewski, S. (1902). Über Flächen des vierdimensionalen Raumes, deren sämtliche Tangentialebenen untereinander gleichwinklig sind, und ihre Beziehung zu ebenen Kurven. Diss. Zürich, Verlag von F. Speidel Zürich.
[Lavrentiev and Chabat, 1980] Lavrentiev, M. and Chabat, B. (1980). Effets hydrodynamiques et modèles mathématiques. Traduit du russe. Mir, Moscou.
[Maruyama, 1950] Maruyama, T. (1950). Some properties on geometry in complex space. Science Journal of the Gakugei Faculty, Tokushima University, 1:31-35.
[Parlett, 1998] Parlett, B. (1998). The Symmetric Eigenvalue Problem. Classics in Applied mathematics, volume 20. SIAM.
[Price, 1991] Price, G. B. (1991). An introduction to multicomplex spaces and functions, volume 140. CRC Press. [Rincon-Camacho, 2015a] Rincon-Camacho, M. M. (2015a). Revisiting the relative efficiency of linear models in statistics and econometrics (in preparation). Technical report, CERFACS.
[Rincon-Camacho, 2015b] Rincon-Camacho, M. M. (2015b). Unpublished manuscript.
[Rincon-Camacho and Latre, 2013] Rincon-Camacho, M. M. and Latre, J.-B. (2013). Project description: Algebraic structures for information processing associated with Dickson algebras. Technical report, CERFACS TR/PA/13/106.
[Scharnhorst, 2001] Scharnhorst, K. (2001). Angles in complex vector spaces. Acta Applicandae Mathematica, 69(1):95-103.
[Segre, 1892] Segre, C. (1892). Le rappresentazioni reali delle forme complesse e gli enti iperalgebrici. Math. Ann, 40:413-467.
[Wang and Chow, 1994] Wang, S.-G. and Chow, S.-C. (1994). Advanced linear models: theory and applications. M. Dekker.
[Wong, 1977] Wong, Y.-C. (1977). Linear geometry in Euclidean 4-space. Number 1. Southeast Asian Mathematical Society.


[^0]:    ${ }^{(1)}$ CERFACS, 42, avenue Gaspard Coriolis, 31057 Toulouse Cedex 1, France. (chatelin@cerfacs.fr) .
    ${ }^{(2)}$ CEREMATH, Université Toulouse 1, 21, Allée de Brienne, 31000 Toulouse, France.

